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Title	ナノ粒子分布制御に基づく熱伝導性ポリマーナノ複合 材料の設計
Author(s)	張,習
Citation	
Issue Date	2020-09
Туре	Thesis or Dissertation
Text version	ETD
URL	http://hdl.handle.net/10119/17007
Rights	
Description	Supervisor:谷池 俊明,先端科学技術研究科,博士



Japan Advanced Institute of Science and Technology

**Doctoral Dissertation** 

# Design of Thermal Conductive Polymer Nanocomposites Based on Nanoparticle Distribution Control

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#### Design of Thermal Conductive Polymer Nanocomposites Based on Nanoparticle Distribution Control

#### Zhang Xi 1720418

Along with the miniaturization and lightweight design of electronic devices, greater demands are constantly raised on materials having high thermal conductivity and low density. Polymer nanocomposites are excellent candidates owing to the general characteristics of polymer materials such as good processability, light weight, high electrical resistivity, and low cost. However, the present polymer nanocomposites usually exhibit insufficient thermal conductivity unless impractical formulations are adopted, and key factors affecting the formation of thermal conductive networks are still unclear. In this thesis, I aimed to study the thermal conductivity. Polymer nanocomposites with continuous thermal conductive networks were designed based on controlling the selective distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles. The influences of filler dispersion, filler migration and phase morphology evolution on the formation of thermal conductive networks were studied, and the relationship of filler distribution and thermal conductivity was investigated. The main research results are as follows:

In **Chapter 2**, I applied a reactor granule technology (RGT) to immiscible polypropylene (PP)/polyolefin elastomer (POE) blends to control the localization of nanoparticles at the interface of PP and POE. The RGT afforded uniform dispersion of in-situ generated Al<sub>2</sub>O<sub>3</sub> nanoparticles in PP, and this guaranteed the migration of nanoparticles to the interface of a co-continuous structure when blended with POE. The selective localization of nanoparticles at the interface was never achieved when preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles were used, and this fact stresses the importance of uniform dispersion in controlling the migration of nanoparticles.

In **Chapter 3**, the relationship among the filler dispersion, phase morphology evolution and the distribution of fillers was revealed by studying the migration of nanoparticles during the annealing process. It was found that the uniform dispersion of nanoparticles decreased the phase domain size to facilitate the successful migration to the interface. Further, the formation of nanoparticle networks decelerated the phase coarsening during annealing. Contrary, the formation of agglomerates and clusters prevented the successful migration of nanoparticles and thus suppression of the phase coarsening became much less effective.

In **Chapter 4**, a continuous segregated structure was designed to improve the thermal conductivity of  $PP/Al_2O_3$  nanocomposites. I achieved the thermal conductivity of 1.07 W/m K at an  $Al_2O_3$  loading of 27.5 vol%. The same strategy was also used in combination with RGT to control the balance between the number and density of thermal conductive paths.

In conclusion, the filler dispersion, phase domain size, and filler distribution in nanocomposites have important cooperation to determine the thermal conductivity of nanocomposites (Fig. 1). RGT and other developed strategies are believed to be promising for designing practical thermal conductive polymer nanocomposites.



Fig. 1 Influence of filler dispersion, phase domain size and filler distribution on thermal conductivity. **Keywords:** Nanocomposites; Thermal conductivity; Network; Immiscible blends; Selective distribution

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### Preface

The present thesis is submitted for the Degree of Doctor of Philosophy at Japan Advanced Institute of Science and Technology, Japan. The thesis is consolidation of results of the research work on the topic "design of thermal conductive polymer nanocomposites based on nanoparticle distribution control" and implemented during October 2017–September 2020 under the supervision of Assoc. Prof. Toshiaki Taniike at Graduate School of Advanced Science and Technology, Japan Advanced Institute of Science and Technology.

**Chapter 1** provides a general introduction of the research field, and accordingly the objective of this thesis. **Chapter 2** introduces selective localization of aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) at interface and its effect on thermal conductivity in polypropylene (PP)/polyolefin elastomer (POE) blends. **Chapter 3** reports cooperative influences of nanoparticle localization and phase coarsening on thermal conductivity of PP/POE blends. **Chapter 4** reports design of continuous segregated PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites and impact of controlled Al<sub>2</sub>O<sub>3</sub> distribution on thermal conductivity. Finally, **Chapter 5** describes the general conclusions of this thesis. The work is original and no part of this thesis has been plagiarized.

Zhang Xi

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### Acknowledgements

I would like to express my sincere gratitude to my supervisor, Assoc. Prof. Toshiaki Taniike, Graduate School of Advanced Science and Technology, Japan Advanced Institute of Science and Technology, who has offered me kind guidance and whose valuable suggestions, incisive comment, constructive criticism, and heartfelt encouragement have contributed greatly to the completion of this thesis. I am grateful to him for his tremendous assistance in developing the framework and revising the draft of this thesis several times. His patience, encouragement, and professional instructions helped me in all the time of research and writing of this thesis.

I would also take this opportunity to thank Prof. Tatsuo Kaneko and Prof. Noriyoshi Matsumi in JAIST for their guidance and encouragement for me. Without their kind support and help, the completion of the present thesis would not have been possible.

I would like to thank the nanocomposite research group of Taniike laboratory for stimulating discussions and constant help. I am also heartily grateful to all the members of Taniike laboratory for their valuable suggestions, cooperation and support.

I would also like to thank the members of my review committee, Assoc. Prof. Kazuaki Matsumura (JAIST), Assoc. Prof. Ken-ichi Shinohara (JAIST), Assoc. Prof. Kosuke Okeyoshi (JAIST), and Assoc. Prof. Yusuke Hiejima (Kanazawa University), who have spent their valuable time to read the thesis, and provided their insightful comments and remarks to enhance the quality of this thesis from various perspectives.

I am also grateful for the scholarship of the Ministry of Education, Culture, Sports, Science and Technology, Japan during my doctoral course studying.

Last but not the least, my gratitude also extends to my family who has been assisting, supporting and caring for me all my life. Special thanks should go to my friends for their utmost care and moral support.

Zhang Xi

Graduate School of Advanced Science and Technology Japan Advanced Institute of Science and Technology May 2020

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# Chapter 1

## **General Introduction**

#### 1.1 Thermal conductive polymer composites

Polymer materials have been widely applied in the fields of electric and electronic devices, energy, biology, medicine, and so on [1-8], because they possess many advantages such as light weight, good processability, low cost, and excellent chemical stability [9–12]. Among them, thermal conductive materials have attracted great attention in recent years owing to the rapidly developing integration, miniaturization and functionalization of electronics, as well as the emergence of new applications such as ultrahigh-voltage electric devices, three-dimensional chip stack architectures, light emitting diodes, flexible electronics, and so on [13–16]. In these applications, efficient thermal dissipation is very important for electronic devices to maintain desired performance, long lifetime and good reliability. For example, personal computers have developed from original big desktops to present compact tablets even with more functions and faster operation, while it inevitably results in the generation of a large amount of heat per volume. If such heat cannot be dissipated, the electronic devices would accelerate the aging and degradation. It was reported that the performance stability of these devices decreases 10% with the increase of temperature by 2 °C and the lifetime decreases to 1/6 if the temperature increases from 25 to 50 °C [17]. Likewise, solving the heat dissipation problem is critical to the electronic information technology, which provides a good opportunity for the design and development of thermal conductive materials.

Thermal conductive polymer materials are generally divided into intrinsic thermal conductive polymers and thermal conductive polymer composites with a thermal conductive inclusion. Intrinsic thermal conductive polymers refer to polymers with high thermal conductivity by themselves. They are usually obtained by regulating the molecular structures during the synthesis and formation process. The thermal conductivity of these polymers is largely determined by the chain morphology. As shown in Fig. 1.1a, a semi-crystalline polymer contains crystalline domains with periodically aligned polymer chains and amorphous domains with randomly entangled polymer chains [18]. The transportation of phonons in the crystalline domain is much more effective than in the amorphous domain; when the former is dominant in a polymer, the vibration mode tends to be localized and the thermal conductivity is thus low. Therefore, enhancing the alignment of polymer chains is expected a promising way to improve the thermal conductivity. According to the molecular dynamic simulations, the thermal conductivity of a single extended polyethylene (PE) chains can be as high as 350 W/m K [19,20]. In literature, a thermal conductivity value of ~ 104 W/m K was achieved by ultra-drawing polyethylene into nanofibers with diameters of  $50 \sim 500$  nm and lengths up to tens of millimeters [21]. Even for polydimethylsiloxane, a polymer with a large barrier for heat flow, a single molecular chain still had a thermal conductivity of  $\sim 7 \text{ W/m}$ K [22]. Virendra et al. developed an electro-polymerization using nanoscale templates to improve the thermal conductivity of polythiophene nanofibers up to ~ 4.4 W/m K and it was more than 20 times than the bulk polymer. The enhancement was attributed to the significant polymer chain orientation [23]. Apart from the fibration to obtain highly aligned polymer chains, a locally ordered structure is also useful for thermal conduction. As shown in Fig. 1.2, Kim et al. utilized strong hydrogen bonding between short and rigid polymer A and long and flexible polymer B to realize the penetration of polymer A within the gyration radius of polymer B. Under the influence of polymer B and strong hydrogen bonding, polymer A tended to form an ordered structure, which greatly reduced the barrier of phonon transportation [24].



Fig. 1.1. Schematic illustration of a polymer: a) the morphology of a polymer consisting of crystalline and amorphous domains, and b) structure of a polymer chain. Reproduced from Ref. [18].



Fig. 1.2. Illustrations of heterogeneous (left) and homogeneous (right) distribution of thermal conductive interchain connections. Reproduced from Ref. [24].

As discussed above, some special polymers equipped with high thermal conductivity have been obtained, while the fabrication of such special polymers is often extremely complicated and expensive. Moreover, the form of the materials is generally a nano-sized solid, which severely limits the applications. Besides, the thermal conductivity of the polymers is largely influenced by other factors such as the chain alignment, chain defects, and interfacial vibration, which would greatly scatter the phonons [25]. Due to these reasons, the study of intrinsic thermal conductive polymers is still in the laboratory exploration stage.

Compared with intrinsic thermal conductive polymers, thermal conductive polymer composites are more attractive and promising candidates for heat dissipation due to the ease, feasibility, and flexibility of the processing. In fact, the thermal conductivity of practical polymers is low, which usually falls in the range of  $0.2 \sim 0.50$  W/m K [7,26–29], due to significant phonon scattering by defects and

grain boundaries in polymer matrices. By adding fillers with high thermal conductivity, the thermal conductivity of polymers can be greatly enhanced while maintaining the good processability of the original polymers. Many thermal conductive fillers can be used for the preparation of thermal conductive composites. Among them, thermal conductive fillers with electrical insulating properties including aluminium oxide (Al<sub>2</sub>O<sub>3</sub>), aluminum nitride (AlN), boron nitride (BN), magnesium oxide (MgO), zinc oxide (ZnO), silicon oxide (SiO<sub>2</sub>), silicon carbide (SiC) have been widely used to fabricate insulating thermal conductive composites [30-37] to satisfy the requirement of high thermal conductivity and electrical insulation. Yao et al. improved the thermal conductivity of glass fiber-reinforced bismaleimide-triazine resin to 1.07 W/m K by adding 70 wt% Al<sub>2</sub>O<sub>3</sub> nanoparticles [30]. Lee et al. obtained epoxy/AlN composites with the thermal conductivity of 3.39 W/m K, 15 times higher than the pure epoxy at the loading of 57 vol% [38]. As such, an extremely high filler loading, typically over 50 vol%, is usually required to achieve the desired thermal conductivity, which is detrimental for the weight, cost and processing. Accordingly, the current research focus in thermal conductive polymer materials is the improvement of the thermal conductivity using a reduced filler loading.

#### 1.2 Applications of thermal conductive polymer composites

Thermal conductive material attracts great attention with the rapidly increasing cooling demand in emerging industries. Polymer composites are very promising for many practical applications due to their advantages, such as light weight, good processability, low cost, and the controllable electrical conduction and insulation. Here, four representative applications of thermal conductive polymer composites are described.

#### **1.2.1** Electronic packaging

The rapid miniaturization and high degree of integration of electronic devices inevitably result in the generation of a large amount of heat in small space. This heat must be efficiently dissipated, otherwise the lifetime and performance of electronic devices will be influenced. The design of electronic packaging is thus very important for heat dissipation and elimination of hot spots in electronic devices. Inorganic thermal conductive materials with good electrical insulating, including metal oxides, nitride ceramics, and other non-metallic materials, were the main electronic packing materials at the early stage. However, they cannot satisfy the demand of the modern electronic packaging technology due to their high price and poor processability. Thermal conductive polymer materials have become a promising alternative and gradually replaced the traditional ceramic electronic packing materials. The application of thermal conductive composites in electronic packaging are mainly focused on automotive electronic control units, single chip packages, and 3D chip packages [13].

#### 1.2.2 LED packaging

Light emitting diode (LED) has higher energy efficiency and longer lifetime than traditional lighting, thus attracting attentions. However, even for high-efficiency LED devices, only  $15 \sim 20$  % of the input energy can be converted to light, and the remaining energy is transferred to heat. If this heat cannot be dissipated from LED devices, the lighting performance, lifetime and energy efficiency can be significantly influenced. Therefore, materials with good thermal conductivity for heat dissipation from the LED inside package to the outside are very important.

Conventionally, the thermal management system of LED devices contains a heat sink and thermal interface materials. The thermal conductivity of commonly used thermal interface materials, including phase-change materials, silicon-based grease and gels, is relatively low. Increasing the thermal conductivity of thermal interface materials or reducing the interfacial thermal resistance can improve the heat dissipation of LED devices. Hu et al. applied CNTs into silicone/metal composites, obtaining a thermal interface material with high thermal conductivity [39]. Aligning CNTs in the direction perpendicularly with the contact surface is regarded as a useful method to further enhance the thermal conductivity of thermal interface materials [40-42].

A new direction to improve the heat dissipation efficiency is directly placing chips on the heat sink. Good electrical insulation of the heat sink is the prerequisite of this method. The commonly used method is filling polymer with thermal conductive ceramic materials. However, a very high filler loading is needed to reach a required thermal conductivity value higher than 10 W/m K. Such a high filler loading brings new problems such as high cost and poor processability. Therefore, achieving high thermal conductivity at a low filler loading becomes a current research target. Besides, the heat dissipation efficiency is also affected by the geometry of the heat sink [43,44].

#### **1.2.3** Aerospace applications

High-performance engineering polymers or their composites are important components in the aerospace area. These materials offer advantages such as excellent mechanical properties, corrosion resistance, lightness, self-lubrication, and economy in fabrication and decoration processes. Generally, aerospace devices should endure hash environments such as high frequency, high pressure and high power factor, to maintain their stability and high performance. In this case, the good heat dissipation must be guaranteed. At present, thermal conductive composites filled with Al<sub>2</sub>O<sub>3</sub> or AlN powder have been widely used to prepare highly thermal conductive composites for packaging and connecting materials in the aerospace area.

#### 1.2.4 Batteries

Thermal management is also very important for batteries especially because that they are more powerful and applied in many devices. More heat can be generated than the dissipated heat when the charge/discharge rate is high or when the batteries are used at high temperature. If too much heat is concentrated in the batteries, the bursting could happen due to the internal pressure buildup. Due to this reason, materials with high thermal conductivity are needed to effectively dissipate the heat inside the batteries, for which thermal conductive composites have been used. For example, Maleki et al. found a hybrid graphite and carbon black filled PVDF composite is a good alternative for the traditional negative electrode in Li-icon cells, due to the high thermal conductivity [45]. Khan et al. fabricated graphene nanoflakes embedded polyacrylonitrile and polymethyl methacrylate nanofiber separators by electrospinning [46]. The thermal conductivity became as high as 8.5 W/m K by the addition of 8 wt% of graphene.

#### **1.3 Mechanisms and models of thermal conductive composites**

#### **1.3.1** Thermal conductive mechanisms

Heat transfer concerns about the generation, use, conversion, and change of heat and is usually realized in the means of thermal conduction, thermal radiation, and thermal conversion. According to the thermal dynamic theory, it associates with the movement, vibration, and rotation of molecules, atoms and electrons [10,47,48]. A mater shows different thermal conductivity at different states, while no matter in what states, the thermal conduction transfers the vibrational energy to the neighboring one and is resulted from the collision and without any motion of mater (Fig. 1.3) [11,49,50].



Fig. 1.3. Heat conduction caused by collision. Reproduced from Ref. [11].

The thermal conduction carrier in a solid mainly includes electrons, photons and phonons. For most polymers, the primary thermal conductivity is dominated by phonons [51]. In the crystalline domains, the thermal conduction comes from the vibration of highly ordered crystalline lattice. When heat is applied, it firstly arrives at the closest surface atom and then flows to the adjacent atom. Such vibration of atoms will be transmitted in a similar way as the propagation of waves through the whole crystal [11]. The quantization of the collective lattice vibration energy is called a phonon [52]. In the amorphous domains, the thermal conduction is based on the independent or non-correlated vibration among the neighboring atoms, where

heat cannot flow as waves, but will diffuse much more slowly. When heat reaches the atoms, it causes disordered vibration and rotation of other atoms around them and then heat is scattered to the adjacent chains [53]. The crystallinity for most polymers is low and the crystals are usually not perfect, thus the thermal conductivity of polymers is limited. Furthermore, inherent defects in polymers including impurities, voids, entanglements, and chain ends can cause considerable phonon scattering, resulting in low thermal conductivity.

The thermal conductive percolation theory and thermal conductive path theory are the main thermal conductive mechanisms of thermal conductive composites. In fact, the percolation theory has been widely used to explain the electrical conductive behavior of electrical conductive composites consisting of electrical insulating polymer matrices and electrical conductive fillers. With the increase of filler loading, materials transfer from insulator to conductor and the electrical conductivity dramatically jumps up over a critical loading, termed the percolation threshold [54–58]. For thermal conductive composites, the percolation theory works only for the system filled with specific fillers such as carbon nanotubes and graphite [59,60]. Many problems appeared when using the percolation theory to explain the thermal conductive behavior. Unlike the sudden increase of electric conductivity at the percolation threshold, the increase of thermal conductivity is much slower. In further increasing the thermal conductive fillers over the percolation threshold, the thermal

conductivity tends to increase linearly rather than keeping a constant value [61]. These problems limit the application of the percolation theory in the field of thermal conductive composites.

Compared with the percolation theory, the thermal conductive path theory is more versatile and widely applied by scholars to explain the thermal conductive behavior of composites. As shown in Fig. 1.4a, individual particles of a thermal conductive filler are isolated in a polymer matrix at a low filler loading. In this case, the polymer matrix constitutes a continuous phase and the filler the dispersed phase, in which filler particles are surrounded by the matrix without any contact or overlap between particles, forming a structure similar to the "sea-island" structure. Therefore, the thermal conductivity of composites under this condition is mainly determined by the thermal conductivity of the polymer matrix. The addition of the filler causes a minor impact on the thermal conductivity. With the increase of the filler loading, the concentration of filler particles in the polymer matrix increases, and some filler clusters are formed (Fig. 1.4b,c). The formed filler clusters contact and overlap with each other and local thermal conductive paths can be thus constructed. When the loading is further increased, thermal conductive paths can gradually connect with each other, and finally develop to a thermal conductive network throughout the whole material (Fig. 1.4d). Under this condition, both of the polymer matrix and filler are continuous and the interpenetrating thermal conductive network significantly improves the thermal conductivity of composites [62].



Fig. 1.4. Thermal conductive composite with different filler loadings. a) fillers are isolated in matrix at extremely low loading, b) some fillers contact with each other when more fillers are added, c) local thermal conductive paths are formed with the further increase of fillers, and d) continuous thermal conductive paths are formed at a high filler loading. Reproduced from Ref. [10].

#### 1.3.2 Thermal conductive models and equations

The thermal conductivity of fillers and polymer matrices, and the interfacial thermal resistance greatly affect thermal conductivity of polymer composites, and it also has a close relationship of the microstructure of polymers as well as the loading, shape, surface performance, dispersion, and distribution of thermal conductive fillers [63]. Therefore, estimating the thermal conductivity of composites by suitable models and equations is very crucial. Based on the thermal conductive path theory, many models and equations were proposed to predict the thermal conductivity of composites [51,64–67].

#### Parallel and series models

Series and parallel models (Fig. 1.5) are two simple and basic models developed to predict the thermal conductivity of polymer composites. The parallel model assumes that the temperature gradient is uniform in composites and the total heat flux is the weight sum of the heat flux in the filler and polymer parts. The series model assumes that the heat flux is uniform in composites and the temperature gradient is the weight sum of the temperature gradient in the filler and polymer parts [13].

According to the parallel model, a polymer matrix and a thermal conductive filler contribute to the overall thermal conductivity independently in proportional to their volume fractions,

$$\lambda = (1 - V_f)\lambda_p + V_f\lambda_f \tag{1-1}.$$

Here,  $\lambda$ ,  $\lambda_p$  and  $\lambda_f$  are the thermal conductivity of the composite, polymer matrix and filler, respectively. V<sub>f</sub> represents the volume fraction of the filler. This model

maximizes the contribution of the filler because it is based on the assumption that the filler particles are perfectly contacted in fully percolated networks. At present, the parallel model has been applied for composites filled with continuous fibers in the direction parallel to the fibers and to estimate the upper bound of the thermal conductivity of composites.



Fig. 1.5. Series and parallel models for thermal conductive composites. Reproduced from Ref. [13]

Compared with the parallel model, the series model assumes that fillers are isolated with each other and no contact between filler particles is present, and thus the model gives the lower boundary. Within this model, the thermal conductivity of composites can be predicted by Eq. (1-2). When thermal conductive filler particles are uniformly distributed in a polymer matrix and no thermal conductive networks are formed in composites, the series model shows high accuracy.

$$\lambda = \frac{\lambda_{\rm p} \lambda_{\rm f}}{\lambda_{\rm f} (1 - V_{\rm f}) + \lambda_{\rm p} V_{\rm f}}$$
(1-2).

It is known that the thermal conductivity of composites usually falls in the region between the parallel and series models. Even so, the series model is generally closer to the experimental value, thus resulting to the development of many other thermal conductive models based on the series model.

#### Hashin-Shtrikman model

The Hashin-Shikman model was developed to predict the thermal conductivity of macroscopically homogeneous and isotropic composites [68]. When the thermal conductivity of the polymer matrix is lower than that of the filler, the thermal conductivity of composites can be estimated by the Hashin-Shtrikman low boundary ( $\lambda^{\text{HS-}}$ ) as,

$$\lambda^{\text{HS-}} = \lambda_p \frac{2\lambda_p + \lambda_f - 2V_f(\lambda_p - \lambda_f)}{2\lambda_p + \lambda_f + V_f(\lambda_p - \lambda_f)}$$
(1-3).

If the polymer matrix processes higher thermal conductivity and surrounds filler particles with lower thermal conductivity, the Hashin-Shtrikman upper boundary  $(\lambda^{HS^+})$  can be used to predict the thermal conductivity of composites,

$$\lambda^{\mathrm{HS}+} = \lambda_{\mathrm{f}} \frac{2\lambda_{\mathrm{f}} + \lambda_{\mathrm{p}} - 2\mathrm{V}_{\mathrm{p}}(\lambda_{\mathrm{f}} - \lambda_{\mathrm{p}})}{2\lambda_{\mathrm{f}} + \lambda_{\mathrm{p}} + \mathrm{V}_{\mathrm{p}}(\lambda_{\mathrm{f}} - \lambda_{\mathrm{p}})}$$
(1-4).

Using the Hashin-Shtrikman low and upper boundaries, the interconnectivity of the thermal conductive part in composites (X<sub>interconnected</sub>) can be evaluated based on,

$$X_{\text{interconnected}} = \frac{\lambda_{\text{measured}} - \lambda^{\text{HS-}}}{\lambda^{\text{HS+}} - \lambda^{\text{HS-}}}$$
(1-5),

which can be used to evaluate the network formation to the ideally interconnected network [69]. The closer to 1 the X<sub>interconnected</sub> is, the better the interconnection of the thermal conductive paths is. The Hashin-Shikman bounds always fall within the parallel-series bounds, regardless of the volume fraction and the thermal conductivity of polymers and fillers, thus they are the best possible boundaries for homogeneous, isotropic binary systems [63,70-73]. However, the Hashin-Shtrikman model assumes a spherical approximation for thermal conductive fillers, is suitable for fillers with high ratio and it not aspect and/or microgeometry/morphology.

#### Bruggeman model

Even though many models have been proved successful to predict the thermal conductivity of composites, they are basically based on the assumption that the thermal conductive fillers are isolated in the polymer matrix. Therefore, these models are not valid for composites with a high filler loading, where some fillers are likely to contact with each other and possibly results in the formation of thermal conductive paths [74]. For such composites, a theoretical model proposed by Bruggeman [75] and modified by Landauer [76] (Eq. 1-6) is more suitable for predicting the thermal conductivity,

$$1 - V_{f} = \frac{\lambda_{f} - \lambda_{c}}{\lambda_{f} - \lambda_{p}} \left(\frac{\lambda_{p}}{\lambda_{c}}\right)^{1/3}$$
(1-6).

This model is based on the differential effective medium (DEM) theory, which assumes the composite can be built up incrementally by introducing infinitesimal changes to a material that already exists. This model uses the mean field approach to capture the interconnection among the fillers that are randomly distributed in the polymer matrix and no limitation is present for the filler concentration [77]. Due to this reason, the Bruggeman model is applicable for the prediction of highly filled composites where filler clusters are formed.

#### <u>Hatta model</u>

The Hatta model is usually applied to predict the thermal conductivity of composites filled with thermal conductive fillers of a platelet shape.

$$\lambda = \left[1 + \frac{V_{f}}{S(1 - V_{f}) + \lambda_{p} / (\lambda_{f} - \lambda_{p})}\right] \lambda_{p}$$
(1-7).

Here, S is related to the measurement direction:  $S = \pi d/4t$  in the plane direction and 1- $\pi d/4t$  in the thickness direction, where d and t are the diameter and thickness of the plates, respectively. This model incorporate the effect of the filler orientation, shape as well as anisotropy in the particle shape, proven to be sufficiently accurate for composites with platelet-shaped fillers, such as graphene and boron nitride [78– 81].

#### Agari model

The Agari model is another suitable theoretical model for polymer composites at a high filler loading by introducing the vertical and parallel conduction mechanism [82–84]. Unlike other models assuming the shape of fillers (usually spherical), Agari model does not define any shape factors of fillers. In addition, the Agari model takes the filler dispersion state into consideration by introducing two factors C<sub>1</sub> and C<sub>2</sub> as,  $\log \lambda = V_f C_2 \log \lambda_f + (1-V_f) \log(C_1 \lambda_p)$  (1-8),

where  $C_1$  is a factor relating to the polymer structure, such as crystallinity and crystal size, and  $C_2$  is a factor relating to the ease of forming thermal conductive paths. The closer to 1  $C_2$  is, the easier for the fillers to form thermal conductive paths. According to the Agari model, the thermal conductivity of composites can vary with the change of filler dispersion state, even if they are consisted with the same components. In other words, based on fitting the experimental data into the Agari model, the dispersion states of fillers in composites can be obtained [85,86].

Apart from the above-mentioned models, other theoretical models such as Hamilton-Crosser model [87], Zhou model [65], Cheng-Vachon model [88], Russell model [89], Baschirow-Selenew model [90], Nielsen-Lewis model [91], and Spinger-Tasi model [92] have also been developed to predict the thermal conductivity of polymer composites. They can reasonably fit the experimental data for specific systems. In general, the thermal conductive models can give an estimate for the thermal conductivity of composites and avoid obviously unnecessary experiments. On the other hand, the models are generally based on some ideal assumptions such as isolated fillers with spherical shapes, perfectly aligned fibers or plates, which are different from the real case. Therefore, more accurate analysis is required and the models need to be further modified [13].

#### 1.4 Main research directions for high thermal conductivity

As discussed in 1.1, preparing thermal conductive polymer composites by adding highly thermal conductive fillers is more feasible than controlling the polymer morphology at an atomic/molecular level. At present, a high filler loading is usually needed to achieve high thermal conductivity, while it compromises other important properties such as mechanical and optical properties. Based on these considerations, improving the thermal conductivity of composites at a low filler loading is a major challenge of the field. The research efforts in this direction is described below, which are subdivided into interface engineering, filler orientation, hybrid fillers, and formation of filler networks as different strategies.

#### 1.4.1 Interfacial engineering

In practice, the enhancement of thermal conductivity in composites is usually unsatisfied because of the high interfacial resistance originated from strong interfacial phonon scattering [93–97]. The aggregation states of polymer matrices and thermal conductive fillers are different, causing inharmonic phonon vibration and interfacial phonon scattering. In addition, defects and voids resulted from poor filler dispersion and interfacial incompatibility between the polymer matrix and fillers hinder the thermal conductivity enhancement. Therefore, engineering the interface between matrices and fillers becomes a trend to fabricate highly thermal conductive composites.

The surface treatment and functionalization of thermal conductive fillers by covalent and non-covalent bonding have been widely applied to enhance the coupling of polymer matrices and fillers. Commonly used covalent modifiers include silane coupling agents, surfactants, grafted polymer chains and etc. [98–101]. According to molecular dynamics simulations, the interfacial thermal resistance of graphene–polymer can be reduced by grafting polymer chains to graphene [102–104]. The improvement of thermal conductivity by surface treatment and functionalization were also proven by many experiments. In our previous research, it was found that a silane coupling agent, (2-phenylethyl)trimethoxysilane, can tailor the interfacial interaction between PP and Al<sub>2</sub>O<sub>3</sub> nanoparticles, leading to the enhancement of thermal conductivity of 0.74 W/m K at 20 wt% compared to 0.41 W/m K for the unmodified nanocomposite at the same loading [105]. Wattanakul et al. improved the thermal conductivity of BN filled epoxy composites from 1.5 W/m K for untreated BN to 2.69 W/m K by admicellar treatment to

enhance the interfacial adhesion [106]. Ding et al. synthesized polyamide-6 (PA6)/graphene (GO) composites by in-situ ring-opening polymerization reaction, in which PA6 chains were covalently grafted onto GO sheets through the "grafting to" strategy. The good interfacial adhesion between PA6 and GO improved the thermal conductivity from 0.196 W/m K for neat PA6 to 0.416 W/m K of PA6/GO composite at a GO loading of 10 wt% [107]. Yu et al. treated Al<sub>2</sub>O<sub>3</sub> nanoparticles with a silane coupling agent and then grafted hyperbranched aromatic polyamide from the modified surface. The resultant Al<sub>2</sub>O<sub>3</sub> nanoparticles were uniformly distributed in an epoxy matrix and the interfacial interaction was improved, leading to enhanced thermal conductivity [108].

Even though the above and other research efforts demonstrated that covalent functionalization of fillers could improve the thermal conductivity of composites, many other results reported that it leads to minor improvement or even could reduce the thermal conductivity of composites. This is because the covalent functionalization may deteriorate the surface structure of fillers, resulting in the loss of the instinct thermal conductivity of fillers [109–111]. In this light, non-covalent functionalization is an alternative approach to enhance the polymer-filler interfacial interaction without reducing the instinct thermal conductivity of fillers. Teng et al. functionalized graphene nanosheets (GNSs) by  $\pi$ - $\pi$  stacking of pyrene having a functional segmented polymer chain in the molecular structure, which resulted in

the significant enhancement of thermal conductivity to 1.91 W/m K,, which was 20% higher than GNS/epoxy and 267% higher than carbon nanotube (CNT)/epoxy [112]. Zhang et al. enhanced the interfacial interaction of CNTs and poly(vinylidene fluoride) (PVDF) by treating CNTs using polyvinylpyrrolidone, resulting in the improvement of thermal conductivity to 0.63 W/m K at the CNT loading of 10 wt%, which was 34% higher than untreated CNT filled PVDF [113]. Wu et al. found that non-covalent functionalization of BN by polydopamine not only improved the thermal conductivity of nanocomposites, but also resulted in outstanding mechanical properties and dimensional stability [114].

The functionalization of thermal conductive fillers can improve the interfacial performance of polymer and fillers. However, the introduced organic modifier is an isolation layer of heat and overmuch modification can reduce the thermal conductivity. In contrast, inorganic coating can not only enhance the interfacial interaction between polymer and fillers, but also further improve the thermal conductivity by employing thermal conductive layer [115]. Besides, thermal conductive materials need to be electrically insulating for many applications. However, many thermal conductive fillers such as metal particles and carbon materials have good electrical conductivity. By coating the surface of electrically conductive fillers with insulating inorganic materials, electrical insulating polymer composites with high thermal conductivity can be obtained. Cui et al. successfully

prepared core-shell structured MWCNT@SiO<sub>2</sub> with a uniform SiO<sub>2</sub> coating on the surface of MWCNTs, and added them into an epoxy matrix [93]. They found the less stiff SiO<sub>2</sub> intermediate shell introduced on the surface of MWCNTs not only reduced the modulus mismatch between the soft epoxy matrix and the stiff MWCNTs, but also enhanced the interaction, resulting in the increase of thermal conductivity by 67% at a low filler loading of 1 wt%. Similar enhancement of thermal conductivity was reported by Chen et al. using a sol-gel method to synthesized SiO<sub>2</sub>-coated silver nanowires [116], and by Choi et al. using a polyvinylpyrrolidone (PVP)-assisted sol-gel reaction to prepare SiO<sub>2</sub>-coated graphite [117]. Al<sub>2</sub>O<sub>3</sub> nanoparticles were also used to form the core-shell structure. Sun et al. utilized supercritical carbon dioxide (scCO<sub>2</sub>) to form Al<sub>2</sub>O<sub>3</sub> nanoparticlecoated graphene nanoplatelets (GNPs) and then used them to prepare epoxy-based composites [118]. The obtained composites exhibited a high thermal conductivity of 1.49 W/m K at a filler loading of 12 wt%. Al<sub>2</sub>O<sub>3</sub>-coated aluminum (Al) nanoparticles were also proposed to obtain electrical insulating materials with high thermal conductivity by Badi et al. [119].

#### 1.4.2 Filler orientation

One dimensional (1D) and two dimensional (2D) fillers such as fibers, CNTs, nanowires, palate-like BN, graphene sheets and graphite flakes usually have anisotropic thermal conductivity [13]. The axial or in-plane thermal conductivity of
these fillers is generally higher than that of the other direction. For example, the inplane thermal conductivity of BN is  $\sim 600$  W/m K, much higher than  $\sim 30$  W/m K in the out-of-plane direction [120]. For composites filled with 1D or 2D thermal conductive fillers, a perfect orientation of fillers in their longitudinal axes can result in great thermal conductivity enhancement in the direction parallel to the orientation direction. The orientation of fillers is also advantageous for the formation of thermal conductive networks; thus, this strategy is very powerful for improving the thermal conductivity of composites at a low filler loading.

Research efforts on the filler orientation include those related to the in-plane orientation and through-plane orientation. For the former, 1D or 2D fillers preferentially align in the in-plane direction of the thermal conductive composite films and form thermal conductive networks in the in-plane direction. In this case, the thermal conductivity of the composites in the in-plane direction is much higher than that in through-plane direction. Zeng et al. prepared BN nanotube (BNNT)/cellulose nanofiber nanocomposites by vacuum filtration (Fig. 1.6 a,b) [121]. They found that BNNTs were uniformly dispersed in the cellulose nanofiber matrix because of the strong interaction of fillers and the matrix. Due to the gravity and vacuum filtration, BNNTs presented an obviously orientated distribution (Fig. 1.6 c,d). The obtained composites had a thermal conductivity value of 21.39 W/m K in the in-plane direction at a BNNTs

loading of 25 wt%. Kumar et al. utilized a simple solution casting method to fabricate high self-aligned large-area reduced graphene oxide (rLGO)/poly (vinylidene fluoride-co-hexafluoropropylene) (PVDF-HFP) composite films [122]. The obtained film revealed ultrahigh in-plane thermal conductivity of ~ 19.5 W/m K at a filler loading of 27.2 wt%. Apart from these methods, other polymer processing approaches such as two-roll milling [123], melt extrusion [124,125], injection molding [126–128], and spinning [129] can be used to prepare in-plane thermal conductive composites.



Fig. 1.6. a) Schematic illustration of the fabrication of the CNF/BNNT nanocomposites. b) Optical images of CNF/BNNT nanocomposites with different BNNT loadings. c) Surface SEM and d) cross-section morphology of CNF/BNNT nanocomposites with 25 wt % of BNNTs. Reproduced from Ref. [121].

Compared with the in-plane thermal conductivity, increasing the out-of-plane thermal conductivity of composites are more important for practical electronic applications. Unfortunately, it is still a challenge and relatively little research was reported. In recent years, with the technical development of aligning fillers by external fields, such as magnetic and electric fields, and the proposal of an ice-templated assembly strategy, the research of increasing the out-of-plane thermal conductivity of composites has made some progress. Yuan et al. modified the surface of hexagonal BN (hBN) with superparamagnetic iron oxide nanoparticles, then aligned fillers by an external magnetic field (Fig. 1.7) [130]. The resultant composites revealed 44.5 % higher out-of-plane thermal conductivity than unaligned composites at a filler loading of 9.14 vol%. Alignment of thermal conductive fillers by magnetic fields for the thermal conductivity enhancement of composites have also been reported by other researchers [131–135].



Fig. 1.7. Illustration and SEM micrographs of epoxy/hBN composites. Reproduced from Ref. [130].

An electric field is also used to align thermal conductive fillers for enhancing the out-of-plane thermal conductivity. Cho et al. used a high DC electric field to form linear assemble hBN nanosheets (Fig. 1.8a) [136]. The oriented structure (Fig. 1.8b) resulted in effective thermal conduction in the out-of-plane direction. Fujihara et al. further advanced this strategy to assemble hBN nanosheets into a filament-like structure with denser population and extended length in a polymer matrix by microscopic molds [137]. The thermal diffusivity increased from 8.67 × 10<sup>-8</sup> m<sup>2</sup>/s for random distribution to  $9.04 \times 10^{-8}$  m<sup>2</sup>/s for the composites with oriented distribution.



Fig. 1.8. Schematic illustration of a) formation of linear assembled hBN under a DC electric field, and b) heat diffusion in composites with different hBN distribution. Reproduced from Ref. [136].

In recent years, an ice-templated assembly strategy was proposed to orient thermal conductive fillers along the ice-growth direction. Zeng et al. reported the fabrication of a 3D boron nitride nanosheet network by an ice-templated strategy and then infiltration (Fig. 1.9) [138]. The resulting composites exhibited a high out-of-plane thermal conductivity of 2.85 W/m K at the filler loading of 9.29 vol%. This strategy is very effective to design the thermal conductive network structure for high out-of-plane thermal conductivity by controlling the crystalline state of ice in the scaffold during the freezing process, thus it has attracted great attention [139–146]. Even though efficient thermal conductive networks can be generated by the ice-templated assembly strategy, the obtained composites are usually limited to epoxy-based composites due to the preformed network structure.



Fig. 1.9. Schematic diagram of the preparation of 3D-BNNS aerogels. Reproduced from Ref. [138].

#### 1.4.3 Hybrid fillers

Thermal conductive networks formed by a single kind of fillers usually remain some undesired voids or defects. The addition of another kind of fillers can fill these voids to further improve the thermal conductivity of composites. The combination of two or more fillers with different sizes or shapes often shows a synergistic effect and can greatly enhance the thermal conductivity by improving filler dispersion, forming bridges between fillers and maximizing the filler packing density. Besides, the utilization of hybrid fillers can significantly reduce the filler loading, thus it can reduce the difficulty of processing. Due to these advantages, hybrid fillers have been widely used to improve the thermal conductivity of composites [24].

For thermal conductive composites, an inevitable problem is the existence of voids between fillers and/or between the polymer matrix and fillers, especially at a high filler loading. They would break the local thermal conductive networks and significantly scatter phonons at the interface, hindering the enhancement of thermal conductivity. At the same time, it is very difficult to fill these voids by adding fillers with the same size due to the space limitation. Therefore, adding fillers with the same shape but different sizes shows advantages for optimizing the packing density of fillers. As shown in Fig. 1.10, Choi et al. hybridized aluminum nitride (AIN) particles and Al<sub>2</sub>O<sub>3</sub> particles with different sizes [147]. Two types of composites were prepared: one filled with large-sized AlN (10 µm) and small-sized Al<sub>2</sub>O<sub>3</sub>

nanoparticles (0.5  $\mu$ m), and the other filled with large-sized Al<sub>2</sub>O<sub>3</sub> (10  $\mu$ m) and small-sized AlN (0.1  $\mu$ m) particles. High thermal conductivity values of 3.402 W/m K and 2. 842 W/m K were achieved for these two types of composites at a total filler loading of 58.4 vol%. The enhancement of thermal conductivity by the hybridization of fillers with different sizes was reported in Al<sub>2</sub>O<sub>3</sub> and BN composites [148,149].



Fig. 1.10 Schematic illustration for the synergistic effect of hybrid fillers with different sizes. Reproduce from Ref. [147].

Hybridization of fillers with different shapes can also improve the thermal conductivity of composites by creating bridges between fillers. To be specific, the combination of 0D, 2D and 3D fillers are usually used to form thermal conductive networks as shown in Fig. 1.11. Generally, fillers with a high aspect ratio can serve as bridges between other fillers in the hybrid system, thus increasing the thermal

conductive paths. For example, Cao et al. reported an obvious synergistic effect between SiC (0D) and MWCNTs (1D) in PVDF/polystyrene (PS) blends, where the thermal conductivity of composites was increased from 0.98 W/m K at a SiC loading of 11.4 vol% to 1.85 W/m K by adding 2.9 vol% MWCNTs [24]. Such improvement of thermal conductivity was attributed to the bridging effect of MWCNTs among SiC nanoparticles. The similar synergistic effect of hybrid fillers was also reported using 1D and 2D fillers such as CNT and GNP [150], and CNT and BN [151–156], as well as the hybridization of 0D and 2D fillers such as Al<sub>2</sub>O<sub>3</sub> and graphene [148], and Al<sub>2</sub>O<sub>3</sub> and BN [157].



Fig. 1.11 Schematic illustration for the synergistic effect of hybrid fillers with different shapes. Reproduce from Ref. [18].

#### 1.4.4 Formation of filler networks

Forming continuous filler networks is very critical to improve the thermal

conductivity of composites. When thermal conductive fillers are uniformly distributed in the polymer matrix, the thermal conductivity cannot be effectively improved due to the high interfacial thermal resistance, and as a consequence, as 60  $\sim$  70 vol% filler loading is needed to form a continuous filler network [13]. However, such a high filler loading leads to the difficulty in processing and high cost. Many approaches have been proposed to construct continuous network at a low filler loading.

A self-assembly approach has been used to construct continuous thermal conductive networks. The thermal conductive networks can be self-constructed by sintering fillers at low temperatures during the polymer curing. Nanoparticles are usually employed in this approach as they tend to sinter at lower temperatures than microparticles [158,159]. Pashayi et al. designed a three-step method, consisting of self-assembly by i) diffusion-controlled aggregation of polyvinylpyrrolidone (PVP) coated nanoparticles, ii) removal of PVP coating from the surface, and iii) sintering of silver nanoparticles in networked structures (Fig. 1.12), to construct dendritic networks of silver in a epoxy resin [160]. By careful self-assembly control and multistep condition design, a high thermal conductive value of 12.05 W/K m at a silver concentration of 10 vol% was achieved, which is up to 300% greater than the state-of-art composites at the similar silver loading.



Fig. 1.12. Three-step formation of tree-shape thermal conductive networks. Reproduce from Ref. [160].

In recent years, self-supported 3D interconnected filler networks, such as 3D carbon foams (Fig. 1.13) and graphite foams have been used to reduce the high interfacial thermal resistance [161–163]. Due to the rigidity of these foams, they can maintain stable continuous networks for phonon transport. Besides, the 3D foams have high porosity and large pore sizes, which allow the introduction of polymer matrix to prepare composites [164,165]. For example, Li et al. incorporated polyamide-6 (PA6) into a graphene foam. The obtained composites exhibited the thermal conductivity of 0.847 W/m K at the graphene loading of only 2 wt% [166]. The enhancement of thermal conductivity of composites by 3D interconnected filler networks can be also found in other reports [161,167–169]. However, 3D foams face many challenges such as complex fabrication, high cost, as well as the limitation of the matrix into thermoset polymers like epoxy.



Fig. 1.13 SEM micrograph of a carbon foam. Reproduce from Ref. [163].

Recently, a segregated structure has been widely applied for reducing the percolation threshold in electrical materials by constructing local electrical conductive networks [119,170]. A segregated structure means a segregated filler dispersion state where the polymer particles are surrounded by fillers. This structure is usually constructed by mechanically mixing fillers and polymer particles, followed by compression molding of the mixture. When thermal conductive fillers are used to fabricate such a structure, the fillers can contact with each other at a low filler loading. The traditional preparation process of a segregated structure is shown in Fig. 1.14 [171]. First, fillers and polymer particles are mixed in a solvent or physically mixed to enhance the absorption of fillers on the surface of polymer particles. The obtained mixture is compression-molded at a temperature around the melting point of the polymer. Hu et al. [171] fabricated Polypropylene (PP)/AIN composites with a segregated structure by grinding PP particles and AlN powder,

and subsequent compression molding. The resulting composites showed more than 20% higher thermal conductivity than that of composites prepared by solution mixing and melt mixing. Other research also indicated the improvement of thermal conductivity by the segregated structure [172,173]. A potential concern of this strategy is imperfection of networks, e.g. physical blending may lead to non-uniform coating of polymer particles by fillers, random packing of particles may remain voids even after the hot-pressing, and so on.



Fig. 1.14. Schematic illustration for the preparation of PP/AlN composites with a segregated structure Reproduce from Ref. [171].

Distribution of fillers in the matrix is a critical factor that determine the formation of thermal conductive networks. Increasing the local concentration of thermal conductive fillers can effectively help the formation of continuous thermal conductive networks, realizing high thermal conductivity at a low loading. The cocontinuous structure in immiscible two-phase polymer blends has been used to enhance the electrical conductivity of polymer composites [174-177]. In this structure, both two phases are continuous and fillers are selectively distributed in one phase (Fig. 1.15) due to the different affinity of the fillers to two phases, which greatly increase the local concentration of fillers in the one phase at a low filler loading. Only recently, this approach has been applied for thermal conductive materials. Cao et al. prepared SiC-filled PS/PVDF blends with a co-continuous structure [25]. They realized the selective distribution of SiC in the PVDF phase by controlling the melting sequence, resulting in the thermal conductivity value of 1.88 W/m K at a filler loading of 23.1 vol%. They further applied this structure for hybrid fillers (CNT and SiC), and obtained the thermal conductivity of 1.85 W/m K with 2.9 vol% CNT and 11.4 vol% SiC [24]. Huang et al. prepared a masterbatch of poly(lactic acid) (PLA)/graphene by solution mixing [178]. Then the masterbatch was solution-mixed with poly(*\varepsilon*-caprolactone) (PCL), during which fillers were driven to the interface of a continuous structure of PLA/PCL blends. Due to the selective localization of fillers at the interface, the thermal conductivity of composites reached 1.28 W/m K at the filler loading of 0.53 vol%. Compared to the above strategies, the utilization of the co-continuous structure to control the selective distribution of thermal conductive fillers is promising to be applied in many systems due to the easy processability and low cost. However, relatively little

research has been conducted at present, and many important questions for controlling the filler distribution are still open and need to be studied in detail.



Fig. 1.15. Schematic illustrations for the selective distribution of a) low- and b) high-aspect ratio fillers [179].

From the above literature review, it can be known that many strategies have been proposed to improve the thermal conductivity of polymer composites. However, most of these strategies need complicated, i.e. impractical processes otherwise the use of prohibitively expensive fillers (remember one of the most important advantages of polymers is its cheapness). Besides, a majority of approaches are only suitable for a specific polymer matrix, such as epoxy resins. These disadvantages significantly limit the potential applications of thermal conductive polymer composites. Therefore, studying the key factors that affect the thermal conductivity, and based on which to optimize the existed strategies for new methods that are more versatile and cheaper is very important.

## 1.4.5 Polyolefin-based composites and reactor granule technology

Polyolefins are considered as one of the most commonly used thermoplastics due to their good environmental compatibility and excellent mechanical properties. However, it is extremely challenging to fabricate performant polyolefin-based nanocomposites, including thermal conductive ones, due to the chemical inertness of polyolefin against inorganic fillers, compared to other polar polymers. Especially for nano-sized fillers, in most of the cases, they cannot disperse in the polyolefin matrix, resulting in the formation of large and compact filler agglomerates. In such a case, it is extremely difficult for fillers to effectively form thermal conductive network. In addition, the existence of these agglomerates inevitably causes significant reduction in other properties such as mechanical properties and transparency [180]. Therefore, it is very important to develop strategies to avoid the formation of filler agglomerates in polyolefin matrices.

The olefin polymerization based on standard heterogeneous catalysts necessarily produces spherical polymer granule with a 3D porous architecture, within which monomers can diffuse and get polymerized to continue the growth [181]. In another words, the porous polymer particles can itself become a polymerization reactor. This technology is called the reactor granule technology (RGT).

For heterogeneous olefin polymerization, the particle growth usually follows the

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"multi-grain" model [182] as shown in Fig. 1.16. According to this model, the catalyst is comprised of multi-grained particles. During the polymerization, the initial catalyst support becomes fragments and dispersed with the growth of polymer. Monomers polymerize on active centers of these catalyst fragments, forming a polymer shell and causing the catalyst grain to further expand progressively [183]. Thus, obtained polymer granule with a porous architecture is called reactor granule. Generally, the obtained reactor granule is immediately melt-mixed into pellets for easy shipping, and the porous structure is not normally utilized except reactor alloying for high-impact PP [184].



Fig. 1.16. Schematic representation of the "multi-grain" model. Reproduced from Ref. [185].

Our research group has successfully applied the pore confinement concept to in-

situ fabricate polyolefin nanocomposites, which involves the impregnation of a molecular precursor of fillers and subsequently conversion into nanoparticles using melt-mixing (Fig. 1.17). Many nanoparticles, including titanium dioxide (TiO<sub>2</sub>), Al<sub>2</sub>O<sub>3</sub>, magnesium hydroxide (Mg(OH)<sub>2</sub>) gold (Au) and silver (Ag), have been synthesized in the pores of polyolefin based on RGT [186–188]. Due to the pore confinement (Fig. 1.17a,b), extremely uniform dispersion (Fig. 17c) was achieved. It was found that the uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles attributed to a higher thermal conductivity at a lower loading in the absence or presence of a compatibilizer (maleic anhydride grafted PP) (Fig. 17d) [186]. For further improvement, the group also melt-mixed impregnated reactor granule with pristine reactor granule to create the distribution of the dispersion [105]. These results indicate the potential of RGT in filler distribution control for polyolefin composites with high thermal conductivity (Fig. 1.19).



Fig. 1.17 Schematic illustrations for the reactor granule technology.



Fig. 1.18. Cross-sectional SEM micrographs of neat PP reactor granule and b) impregnated with Al(O*i*Pr)<sub>3</sub>. c) TEM micrograph of the PP/Al<sub>2</sub>O<sub>3</sub> nanocomposite prepared by RGT at the Al<sub>2</sub>O<sub>3</sub> loading of 20 wt%. d) Thermal conductivity of three types of PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites. Reproduced from Ref. [186].



Fig. 1.19. Distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles in PP for samples prepared by meltmixing impregnated reactor granule with pristine reactor granule a,b,c) and standard RGT d,e,f). g) Thermal conductivity of two types of PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites. Reproduced from Ref. [105].

# 1.5 Aim of thesis

The rapid development of electronic integration technology results in the greatly increased packing density and remarkably reduced volume of electronic devices and logic circuits. This necessarily accompanies a large amount of heat generated and accumulated in a limited space. Efficient dissipation and elimination of this heat become a key issue to guarantee the reliable operation of relevant devices. Under such circumstances, thermal conductive composites have emerged as attractive materials in many fields, such as micro-electronic devices, LED, aerospace, batteries, and etc. These materials offer advantages such as light weight, good processability, low cost, excellent chemical stability, and electrical insulation, thus attracting great attention in recent years. To obtain polymer composites with desired thermal conductivity, many strategies have been proposed. However, these methods generally aim at improving the thermal conductivity of specified systems, and usually require complex processing as well as very expensive thermal conductive fillers, which greatly increase the fabrication cost and significantly limit the application of obtained materials. In this light, it is very important to study the thermal conductive composites based on more commonly used polymer matrices and thermal conductive fillers. Therefore, this thesis focuses on polyolefin as the matrix, because it is one of the most commonly used thermoplastics with balanced properties and price. Spherical Al<sub>2</sub>O<sub>3</sub> nanoparticles were used as the thermal conductive fillers. Even though, the thermal conductivity of Al<sub>2</sub>O<sub>3</sub> is only moderate (~ 30 W/m K), it has advantages of abundance, electrical insulation and low cost [108,189]. Besides, the incorporation of spherical particles guarantees isotropic heat release in the resultant composite materials [190].

In this thesis, I aim to study the thermal conductivity of polyolefin composites filled with Al<sub>2</sub>O<sub>3</sub> nanoparticles, and to clarify the key factors that affect the thermal conductivity. For this, I designed polymer nanocomposites with continuous thermal conductive networks by controlling the selective localization of nanoparticles. The influence of filler dispersion, filler migration and phase morphology evolution on the formation of thermal conductive networks was studied, and the relationship of filler distribution and thermal conductivity was investigated. My research is mainly comprised of the following three chapters:

In Chapter 2, the reactor granule technology (RGT) was applied to immiscible polypropylene (PP)/polyolefin elastomer (POE) blends to effectively control the localization of nanoparticles at the interface. This study aims to reveal the importance of filler dispersion for controlling the distribution and realize the selective localization at the interface.

In Chapter 3, PP/POE/Al<sub>2</sub>O<sub>3</sub> composites were annealed with different durations and the filler migration and phase morphology evolution were studied by electron microscopy. The relationship among nanoparticle migration, phase morphology evolution, and the thermal conductivity were investigated in detail.

In Chapter 4, a continuous segregated structure was designed to improve the thermal conductivity of PP-based nanocomposites. This proposal was based on the selective distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the POE phase and the compression of a porous scaffold by eating POE of PP/POE blends with a co-continuous structure. The formation of this structure achieved great improvement of thermal conductivity.

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# Chapter 2

# Selective Localization of Aluminum Oxide at Interface and

# Its Effect on Thermal Conductivity in Polypropylene/

**Polyolefin Elastomer Blends** 

### Abstract

A reactor granule technology (RGT) was applied to immiscible polypropylene (PP)/polyolefin elastomer (POE) blends to effectively control the localization of nanoparticles at the interface. The RGT afforded uniform dispersion of in-situ generated aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) nanoparticles in PP, and this guaranteed the migration of nanoparticles to the interface of a co-continuous structure when blended with POE. The selective localization of nanoparticles at the interface was confirmed by transmission electron microscope (TEM), scanning electron microscope (SEM), elemental mapping (EDX), and thermal gravimetric analysis (TGA). Such localization was never achieved when preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles were used, and this fact stresses the importance of uniform dispersion in controlling the migration of nanoparticles. The migration was also affected by the viscosity of POE: POE having higher viscosity resulted in a smaller domain size, and thereby more nanoparticles could migrate to the interface without agglomeration. The impacts of the selective localization at the interface and the phase domain size on the thermal conductivity of the resultant nanocomposites were studied and discussed. Keywords: Thermal conductivity; Nanocomposite; Polypropylene; Immiscible blends; Interface

## 2.1. Introduction

Along with the miniaturization and lightweight design of electronic device components, greater demands are constantly raised on materials having high thermal conductivity and low density [1–4]. Polymeric materials are excellent candidates owing to their general characteristics such as good processability, light weight, low water absorption, corrosion resistance, high electrical resistivity, and most importantly, low cost [5–12]. However, a majority of these polymeric materials possess poor thermal conductivity, which is usually restricted in the range of 0.10–0.25 W/m K [13]. As a consequence, electrically insulating nanoparticles with high thermal conductivity such as boron nitride (BN) [14], aluminum nitride (AIN) [15], silicon carbide (SiC) [16], and aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) [17] are incorporated into polymers for applications that require both high thermal conductivity and electrical insulating properties.

Among these thermal conductive fillers, spherical Al<sub>2</sub>O<sub>3</sub> particles have been widely used due the abundance [18–20]. The incorporation of spherical particles guarantees isotropic heat release in the resultant composite materials [21], but at the same time, necessitates a relatively high filler loading to construct a well-developed thermal conductive network [22,23]. A tradeoff between the thermal conductivity improvement and problems arising from a high loading such as complex processing, high density and severe particle agglomeration must be dealt with in the materials design. Our research group has disclosed a novel reactor granule technology (RGT) to achieve uniform dispersion of nanoparticles in polyolefin-based nanocomposites and such uniform dispersion can be also realized even at an elevated loading. It is based on the impregnation and confinement of metal alkoxides in the pores of polymer reactor granule and subsequent chemical conversion of precursors into nanoparticles during a melt-mixing process [24–26]. In the past research, the group fabricated PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites based on this technology. Owing to the extremely uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles, the improvement of thermal conductivity was achieved at a much lower loading compared with samples prepared by melt-mixing preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles in the absence or presence of a compatibilizer (maleic anhydride grafted PP) [24]. For further improvement, the group also melt-mixed impregnated reactor granule with pristine reactor granule to create the distribution of the dispersion. In detail, Al<sub>2</sub>O<sub>3</sub>-rich and polymer-rich domains were created at around 10 µm scale, in which the filler-rich domains offered a thermal conductive pathway and thus the thermal conductivity was improved to some extent [27]. These results encouraged me to further improve the thermal conductivity by tuning the selective distribution of nanoparticles.

An effective strategy to control the distribution of nanoparticles is the employment of immiscible polymer blends with a co-continuous structure. This strategy has been widely applied in designing electrically conductive polymer composites to improve the electrical conductivity by the selective localization of nanoparticles in one phase [28–35], or, best of all, at the interface [36–38]. Many methods have been reported to achieve the selective localization of carbon black (CB) nanoparticles at the interface. For example, Gubbels et al. reported that compression molding of dry mixed polyethylene (PE), polystyrene (PS), and CB powder resulted in the localization of CB nanoparticles at the interface of PE and PS domains [39]. They also reported in the same paper that by first distributing CB nanoparticles in the less favored phase, the nanoparticles were thermodynamically driven to the interface at proper processing conditions [39]. Al-Saleh et al. proposed another strategy to utilize a compatibilizer for the selective localization of CB nanoparticles at the interface, where a styrene-butadiene-styrene (SBS) tri-block copolymer as a compatibilizer preferentially distributed at the interface of PP and PS domains, and CB nanoparticles were carried to the interface through the attractive interaction with SBS [36].

Recently, this co-continuous structure has also been exploited to design thermal conductive materials. For example, Li et al. introduced BN into poly(3-hydroxylbutyrate)/poly(ethylene oxide) blends with a co-continuous structure. A network of BN was formed in the poly(3-hydroxylbutyrate) phase [40]. Cao et al. achieved high thermal conductivity by controlling the distribution of multi-walled carbon nanotubes and SiC nanoparticles in the poly(vinylidene fluoride) (PVDF)

phase in PVDF/PS blends [41]. Cao et al. reported the selective localization of BN in the polyethylene rubber phase of polyethylene rubber/ethylene propylene diene rubber blends, which yielded high thermal conductivity [42]. All these recent studies revealed the utilization of a co-continuous structure as an effective means for the formation of a thermal conductive network at a reduced filler loading. However, the localization has been attained in one phase, not at the interface, still requiring a relatively high loading (e.g. 10 vol%) [43]. Moreover, the above-mentioned strategies for CB nanoparticles are not directly applicable to Al<sub>2</sub>O<sub>3</sub> nanoparticles because of the extremely high surface tension of Al<sub>2</sub>O<sub>3</sub> (114 mJ/m<sup>2</sup> vs. 55 mJ/m<sup>2</sup> for CB) [44,45]. They are inherently not dispersible in hydrophobic polymer matrices and tend to form agglomerates especially at a high loading [19]. In this light, thermodynamically driving Al<sub>2</sub>O<sub>3</sub> nanoparticles to the interface is very challenging because the agglomeration tendency strongly hinders the migration of individual nanoparticles to the interface.

In this chapter, a novel strategy was proposed based on the RGT and immiscible polymer blends with a co-continuous structure for the first time. The RGT can prevent the agglomeration of in-situ generated Al<sub>2</sub>O<sub>3</sub> nanoparticles and achieve uniform dispersion, which helps the nanoparticles migrate to the interface of immiscible polymer blends. Reactor granule of polypropylene (PP) was preliminarily impregnated with aluminum alkoxide and then melt-mixed with four kinds of POE (used as an assistant phase) with different viscosity. The impacts of dispersion and phase domain size on the migration of nanoparticles were studied.

#### 2.2. Experimental section

#### 2.2.1. Materials

PP reactor granule ( $M_w = 2.6 \times 10^5$ ,  $M_w/M_n = 5.7$ , mmmm = 98 mol%) was synthesized by propylene polymerization using a Mg(OEt)<sub>2</sub>-based Ziegler-Natta catalyst. The median size  $(D_{50})$  and pore volume of the granule were determined as 637 µm and 0.56 mL/g based on laser scattering in ethanol (HORIBA partica, LA-950V2) and mercury porosimetry (Shimazu, Autopore IV 9505), respectively. Aluminium isopropoxide (Al(OiPr)<sub>3</sub>) was supplied by Sigma-Aldrich and used as a precursor. n-Octdecyl-3-(3',5'-di-t-butyl-4'hydroxyphenyl)-propionate bis(1,2,2,6,6-pentamethyl-4-(AO-50) and piperidyl) sebacate (LA-77) were donated by ADEKA. They were used as an anti-oxidant stabilizer and a catalyst for the sol-gel reaction, respectively [46]. Four kinds of POE having different viscosity and octene contents (termed POE1-4) were employed in this study. Table 2.1 summarizes their physical parameters, which are supplied by Dow Chemical Co. POE1 and POE2 have similarly high viscosities compared with those of POE3 and POE4, while POE1 and POE4 have octene contents nearly two times greater than the other two. Such selection was based on the consideration that different viscosity and octene contents can affect the phase domain size and thereby influence the final localization of  $Al_2O_3$  nanoparticles. Besides, the octene content has an impact on the specific heat capacity (Table 2.1) and this can also influence the thermal conductivity of obtained nanocomposites. Preformed  $Al_2O_3$ nanoparticles having the average diameter of 15 nm (Alu, Nippon Aerosil) were used to prepare reference samples.

	POE1	POE2	POE3	POE4
MI (g/10 min) <sup>a</sup>	1	1	30	30
Mooney viscosity (MU) <sup>b</sup>	23	18	<5	<5
Octene content (wt%)	38	20	22	40
Density (g/cm <sup>3</sup> )	0.870	0.902	0.902	0.870
Specific heat capacity (J/g K) <sup>c</sup>	2.81±0.01	2.33±0.01	2.51±0.02	2.83±0.03

Table 2.1. Physical parameters of POE.

<sup>a</sup> Measured at 2.16 kg and 190 °C based on ASTM D1238.

<sup>b</sup> ML 1+4 at 121 °C acquired based on ASTM D1646.

<sup>c</sup> Measured by DSC.

#### 2.2.2. Sample preparation

PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites were obtained based on the following procedure. First, 10 g of PP reactor granule was impregnated with 20 mL of a toluene solution of Al(OiPr)<sub>3</sub> at 50 °C under nitrogen for 12 h. The amount of the Al(OiPr)<sub>3</sub> precursor was determined so as to obtain the  $Al_2O_3$  loading of 20 wt% (= 5.5 vol%) in resultant PP/Al<sub>2</sub>O<sub>3</sub> samples under the assumption of full conversion of the metal alkoxide to the oxide. After the impregnation and solvent removal in vacuo, the impregnated granule was subjected to 50 °C and 100% relative humidity for 24 h to promote the hydrolysis and solidification of Al(OiPr)<sub>3</sub> in the pores of PP [24]. Subsequently, it was melt-mixed with POE at 190 °C and 100 rpm for 10 min. The weight ratio of PP and POE was kept at 1/1 wt/wt to ascertain the formation of a co-continuous structure (Fig. 2.1), and this resulted in the final Al<sub>2</sub>O<sub>3</sub> loading of 11.1 wt% (corresponding to 2.8 vol%). The product was finally hot-pressed into films at 190 °C and 10 MPa for 5 min, followed by quenching at 0 °C for 5 min. The prepared samples were denoted as PP/POEx/Al<sub>2</sub>O<sub>3</sub>-RGT, where x represents the type of POE employed (x = 1-4).

The following reference samples were also prepared.

i) PP/Al<sub>2</sub>O<sub>3</sub>-RGT-*y*: Impregnated PP granule was melt-mixed and hot-pressed in the absence of POE, where *y* represents the volume fraction of Al<sub>2</sub>O<sub>3</sub> in the nanocomposites (y = 2.5 or 5.5 vol%).

ii) PP/POE*x*: Pristine PP was melt-mixed with POE at 1/1 wt/wt, where *x* represents the type of POE.

iii) PP/POEx/Al<sub>2</sub>O<sub>3</sub>-NP: Preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles were directly melt-mixed with pristine PP, followed by blending with POE at the PP/POE weight ratio of 1/1 wt/wt. The final Al<sub>2</sub>O<sub>3</sub> loading was fixed at 2.8 vol%.

The conditions for the impregnation, melt mixing and hot pressing were totally the same as those employed for the preparation of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT samples.



Fig. 2.1. SEM micrographs of etched surfaces of PP/POE blends with different kinds of POE. a,a') POE1, b,b') POE2, c,c') POE3, d,d') POE4. Note that the scale bar is applied to the individual entire row. The formation of a co-continuous structure was confirmed by selectively extracting the POE phase with heptane prior to the SEM measurements.

### 2.2.3. Characterization

The progress of the reaction and the formation of Al<sub>2</sub>O<sub>3</sub> were qualitatively confirmed by recording attenuated total reflection infrared spectra (ATR-IR) (FT-IR, 100, Perkin Elmer) with a resolution of 4  $cm^{-1}$  and 24 scans. The actual Al<sub>2</sub>O<sub>3</sub> contents of prepared samples were measured by thermal gravimetric analysis (TGA, Thermo plus evo, Rigaku). A sample was heated from 25 °C to 600 °C in air at a heating rate of 10 °C/min. The residual inorganic content at 600 °C was regarded as the Al<sub>2</sub>O<sub>3</sub> content. In order to investigate the localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles in a specific phase, PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT samples were etched by heptane to remove POE at 80 °C for 24 h. After rinsing with heptane, ethanol, and deionized water, the samples were subjected to the same TGA measurement. The dispersion and distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the matrix was monitored by a transmission electron microscope (TEM, Hitachi H-7100) operated at an acceleration voltage of 100 kV. TEM specimens with the thickness of 100 nm were prepared by an ultramicrotome (Reichent Ultracut FCS, Reica) equipped with a diamond knife (Diatome). A scanning electron microscope (SEM, Hitachi S-4100) with an acceleration voltage of 20 kV was used to observe the phase morphology of the PP/POE blends. The samples were cryo-fractured in liquid nitrogen and then etched according to the above-mentioned method. The samples were coated with platinum and palladium prior to the measurement. Elemental mapping was also performed

using an energy dispersive X-ray analyzer (EDX, Hitachi TM 3030Plus). The thermal diffusivity ( $\alpha$ ) of samples was measured by a temperature wave analyzer (ai-Phase mobile 1u/2, Hitachi High-Tech Science) [47]. A film sample was sandwiched between the heater and the sensor plates, and the distance between the two plates was adjusted to create a firm contact with both surfaces of the sample. The applied voltage was set to 1.4 V and the phase delay in the temperature wave was measured at eight frequency values within a range of 0.2–2 Hz. The thermal conductivity ( $\lambda$ ) was derived from

$$\lambda = \alpha C_p \rho \tag{2-1},$$

where  $C_p$  is the specific heat capacity and  $\rho$  is the density of a sample. The specific heat capacity at room temperature was determined using a differential scanning calorimeter (DSC, Mettler Toledo DSC-822). The temperature was swept in the range of 0–50 °C at 10 °C/min under nitrogen flow of 200 mL/min. The density was measured by an electronic density meter (SANSYO DME-220).

#### 2.3. Results and discussion

#### 2.3.1. Formation and dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles

ATR-IR spectra was acquired for a  $PP/Al_2O_3$  nanocomposite film to investigate the chemical conversion of the precursor (Fig. 2.2a). Compared to the spectrum of neat PP, the  $PP/Al_2O_3$ -RGT sample shows a new broad band at around 500 cm<sup>-1</sup>, which corresponds to the Al-O-Al vibration of in-situ formed Al<sub>2</sub>O<sub>3</sub> [48]. TGA was employed to verify the actual Al<sub>2</sub>O<sub>3</sub> content (Fig. 2-2b). Based on the residual weight at 600 °C, the actual Al<sub>2</sub>O<sub>3</sub> loading was derived as 20.1 wt%, which was consistent with the theoretical loading (i.e. no evaporation or leaching loss of the precursor).



Fig. 2.2. a) ATR-IR spectra and b) TGA result for a PP/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposite having the theoretical Al<sub>2</sub>O<sub>3</sub> loading of 20 wt%.

The RGT was adopted to realize the uniform dispersion of nanoparticles in the PP matrix. TEM was employed to confirm this expectation. As shown in Fig. 2.3a,b, Al<sub>2</sub>O<sub>3</sub> nanoparticles were uniformly dispersed in the matrix and agglomerates were hardly observed. To further study the dispersion of nanoparticles, a particle size distribution profile was acquired by analyzing TEM images over 150 particles using ImageJ software. The result is shown in the inset of Fig. 2.3b. It can be seen that the size of  $Al_2O_3$  nanoparticles mostly fell in the range of 60–100 nm (the median size ( $D_{50}$ ) = 74.8 nm). Thus, it was confirmed that the RGT generated uniformly dispersed  $Al_2O_3$ nanoparticles in the PP matrix, and this was achieved at a relatively high loading (20 wt%) without adding any dispersant [24].



Fig. 2.3. TEM micrographs of PP/Al<sub>2</sub>O<sub>3</sub>-RGT (20 wt%). The inset shows the particle size distribution of Al<sub>2</sub>O<sub>3</sub>.

#### 2.3.2. Selective localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles in PP/POE blends

To explore the distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles after blending POE, TEM images of the ternary materials were acquired. From Fig. 2.4, the existence of two different phases was clearly observed for all the nanocomposites, while the distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles depended on the viscosity of employed POE.

When POE1 and POE2 having higher viscosity were employed, many Al<sub>2</sub>O<sub>3</sub> nanoparticles can be found at the interface between the PP and POE phases (Fig. 2.4a,b). As was confirmed in **2.3.1**, Al<sub>2</sub>O<sub>3</sub> nanoparticles were uniformly dispersed in the PP matrix, so that the localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles indicated migration of these nanoparticles during the melt processing in the presence of POE. Contrary, the utilization of POE3 and POE4 with lower viscosity caused the confinement of most of Al<sub>2</sub>O<sub>3</sub> nanoparticles in one phase, where particle agglomeration became more evident than the localization at the interface (Fig. 2.4c,d). The different localization behaviours of Al<sub>2</sub>O<sub>3</sub> nanoparticles suggested that the viscosity of POE significantly affected the migration process, and the reason for it will be discussed later.



Fig. 2.4. TEM micrographs of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites with different kinds of POE. a,a') POE1, b,b') POE2, c,c') POE3, d,d') POE4. Note that the scale bar is applied to the individual entire row.

An important role of the viscosity of POE was also confirmed by elemental mapping results for cross-sectioned samples (Fig. 2.5). The distribution of blue and dark regions figures out Al-rich and Al-poor domains, respectively. In the case of POE1 and POE2 (higher viscosity), Al was found concentratedly distributed in specific areas (Fig. 2.5a,b). In contrast, big blue spots were found for POE3 and POE4 (Fig. 2.5c,d), corresponding to the agglomerated Al<sub>2</sub>O<sub>3</sub> nanoparticles.



Fig. 2.5. EDX elemental mapping for Al in PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites with different kinds of POE. a) POE1, b) POE2, c) POE3, d) POE4. The dotted lines figure out Al-rich domains.

In order to identify the phase in which Al<sub>2</sub>O<sub>3</sub> nanoparticles preferred to be distributed in, samples were etched with heptane to selectively extract the POE phase. Then, TGA was used to determine the  $Al_2O_3$  content in the remaining PP phase or at the interface (Fig. 2.6). Before etching the POE phase, the  $Al_2O_3$  content was in the range of 10.2–10.6 wt% (Fig. 2.7), similar to the predicted value (11.1 wt%). After the selective extraction of the POE phase, the  $Al_2O_3$  content was found to be around 20 wt% irrespective of the POE type, and the value was identical to that before adding POE. This fact indicates that most of  $Al_2O_3$  nanoparticles were selectively localized in the PP phase and/or at the interface. Note that the  $Al_2O_3$  nanoparticles were not dispersible in heptane and therefore, the nanoparticles must remain unextracted even though they were localized at the interface.



Fig. 2.6. TGA curves for PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites after selective extraction of the POE phase. The inset summarizes the Al<sub>2</sub>O<sub>3</sub> content derived from the residual weight at 600 °C.



Fig. 2.7. TGA curves for PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites. The inset summarizes the Al<sub>2</sub>O<sub>3</sub> content derived from the residual weight at 600 °C. The Al<sub>2</sub>O<sub>3</sub> content was found to be consistent with the predicted value for all the samples (11.1 wt%).

Cross-sectional SEM images of the etched samples were also acquired (Fig. 2.8). A layer of Al<sub>2</sub>O<sub>3</sub> nanoparticles could be observed at the etched surface of the PP phase in PP/POE1/Al<sub>2</sub>O<sub>3</sub>-RGT (Fig. 2.8a). This corresponds to the localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles at the interface, where the distance between nanoparticles was greatly shortened and many nanoparticles were apparently contacted with each other to form a network. The nanoparticle localization at the interface was similar in PP/POE2/Al<sub>2</sub>O<sub>3</sub> (Fig. 2.8b). In consistent, agglomeration was observed in PP/POE3,4/Al<sub>2</sub>O<sub>3</sub>-RGT (marked in dotted circles in Fig. 2.8c,d), and relatively fewer nanoparticles were found

at the interface. On the basis of these TEM, EDX and TGA results, it was concluded that most of  $Al_2O_3$  nanoparticles were localized at the interface for POE1 and POE2 (higher viscosity), while for POE3 and POE4 (lower viscosity),  $Al_2O_3$  tended to be distributed in the PP phase. The localization at the interface was similarly observed from 0.6 vol% to 4.7 vol% in PP/POE1/Al\_2O\_3-RGT nanocomposites (Fig. 2.9).



Fig. 2.8. SEM micrographs of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites with different kinds of POE after selective extraction of the POE phase. a) POE1,b) POE2, c) POE3, d) POE4. Note that the scale bar is applied to all the samples.



Fig. 2.9. TEM (upper) and SEM (lower) micrographs of PP/POE1/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites with different Al<sub>2</sub>O<sub>3</sub> loadings. a,a') 0.6 vol%, b,b') 1.3 vol%, c,c')
4.7 vol%. Note that the scale bar is applied to the individual entire row.

To clarify the reason why the localization of  $Al_2O_3$  nanoparticles was dominated by the viscosity of POE, the phase morphology was observed by SEM and the results are summarized in Fig. 2.10a–d. The POE phase was preliminarily etched by heptane to give a clear observation. A co-continuous structure was formed in all the samples regardless of the difference in the viscosity and chemical composition of POE. The phase size for PP/POE1/Al<sub>2</sub>O<sub>3</sub>-RGT and PP/POE2/Al<sub>2</sub>O<sub>3</sub>-RGT was determined to be 1.8 µm and 2.6 µm, respectively, which were much smaller than 9.9 µm and 8.6 µm for POE3 and POE4 (Fig. 2.10e). In general, a large interfacial area in polymer blends leads to large interfacial free energy. The phase morphology that is created under an external shear is not stable and it tends to translate to the state with lower interfacial free energy. A direct consequence of this is the coalescence of phases, leading to the increase of the phase size. In the PP/POE3,4 blends, the mobility of polymer chains was much greater so that the POE phase could coalesce and form larger phase morphologies.

According to all the above observations, the morphology of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites is schematically illustrated in Fig. 2.11. When the phase size was large (PP/POE3,4/Al<sub>2</sub>O<sub>3</sub>-RGT), Al<sub>2</sub>O<sub>3</sub> nanoparticles had to migrate farther for arriving at the interface, i.e. they had more chances to encounter with each other and form agglomerates. The agglomeration in turn hindered the further migration of nanoparticles. Thus, fewer Al<sub>2</sub>O<sub>3</sub> nanoparticles could finally arrive at the interface. On the other hand, when the phase size was small (POE1 and POE2), many nanoparticles could migrate to the interface and get stabilized before encountering with each other in the PP phase. The selective localization of nanoparticles at the interface increased their concentration at these areas so that many particles could directly contact with each other to form thermal conductive networks. Besides, a smaller phase size could form more evenly distributed networks in the nanocomposites, and this must be profitable for avoiding the creation of a heterogeneous temperature profile upon heat releasing.



Fig. 2.10. Phase morphology of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites with different kinds of POE after selective extraction of the POE phase. a) POE1,
b) POE2, c) POE3, d) POE4. Note that the scale bar is applied to all the samples. e) Measured phase size.



Fig. 2.11. Schematic illustration of the morphology of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites. The blue domains represent the PP phase, the yellow domains the POE phase, and the black circles Al<sub>2</sub>O<sub>3</sub> nanoparticles.
Here, the role of the polymer viscosity in achieving the localization of  $Al_2O_3$  nanoparticles is discussed from a theoretical viewpoint. It is considered that the localization of nanoparticles is determined by the competition between the speeds of nanoparticle migration and phase coarsening, *e.g.* the successful localization of nanoparticles at the interface can be realized only when the nanoparticle migration is faster than the phase coarsening. The self-diffusion of nanoparticles in a phase that they are initially dispersed in is expressed by the Stokes-Einstein equation for the Brownian motion,

$$D_0 = \frac{k_B T}{6\pi \eta_1 r} \tag{2-2},$$

where  $D_0$  corresponds to the diffusion coefficient of a spherical nanoparticle of radius *r* in a fluid with viscosity  $\eta_1$  (in this study, PP) at temperature *T*, and  $k_B$  is the Boltzmann constant. The time ( $t_D$ ) for the nanoparticle to migrate a distance *r* equal to its radius can be given as Eq. (2-3) [49],

$$t_D = \frac{r^2}{6D_0} = \frac{\pi \eta_1 r^3}{k_B T}$$
(2-3).

The rate of phase coarsening (k) in an immiscible blend follows the equation proposed by Veenstra et al. [50],

$$k = \frac{dR}{dt} = c \frac{\sigma}{\eta_b}$$
(2-4),

where dR/dt describes the rate of the enlargement of the average thickness (*R*) along the time, and  $\sigma$  and  $\eta_b$  are the interfacial tension and the viscosity of the polymer blend respectively. The dimensionless constant *c* is found to be 0.07 [48]. Thus, the time ( $t_C$ ) required for the phase domain size to grow a thickness of r can be obtained as Eq. (2-5).

$$t_C = \frac{r\eta_b}{c\sigma} \tag{2-5}$$

The successful localization of nanoparticles at the interface is expected only when the nanoparticle migration is faster than the phase coarsening, namely  $t_D < t_C$ . Hence the following relation should be satisfied,

$$\frac{\pi\eta_1 r^2}{k_B T} < \frac{\eta_b}{c\sigma}$$
(2-6).

According to Utracki et al. [51] the viscosity of a polymer blend consisting of polymer 1 and polymer 2 at 50/50 vol%/vol% is expressed by

$$\eta_b = \sqrt{\eta_1 \eta_2} \tag{2-7},$$

where  $\eta_1$  and  $\eta_2$  are the viscosity of polymer 1 and polymer 2 at temperature *T*. By substituting Inequality (2-7) into Eq. (2-6), the relationship is re-written as,

$$\sqrt{\frac{\eta_2}{\eta_1}} > \frac{c\pi\sigma r^2}{k_B T}$$
(2-8),

where the left side of the inequality corresponds to the square root of the viscosity ratio of the two polymers. Inequality (2-8) dictates that a higher viscosity of polymer 2 (POE) is more advantageous for the successful localization of nanoparticles at the interface, in agreement with our experimental observation. Meanwhile, a smaller interfacial tension and smaller particle size are other important factors to achieve the localization. The nanoparticle localization is also affected by the temperature dependence of the viscosity ratio (exponential) and 1/T of the right side. For example, when prepared at 190 °C, the successful localization was confirmed for POE1 (and POE2). But, this is not the case when prepared at 230 °C (Fig. 2.12) because the viscosity ratio drops more sharply than 1/T.



Fig. 2.12. TEM micrograph of PP/POE1/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposite prepared at 230 °C.

The preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles were also used to prepare PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP nanocomposites as reference. From Fig. 2.13, poor dispersion of the nanoparticles was evident in all the samples. Such phenomenon is quite common for polyolefin-based nanocomposites because the surface energy of nanoparticles is significantly higher than that of the polymer matrix [52]. More importantly, the Al<sub>2</sub>O<sub>3</sub> nanoparticles were dominantly localized in the PP phase, while only a minor fraction reached the interface, almost regardless of the viscosity of POE. This result suggested the

initially uniform dispersion as another key factor to realize the localization of the nanoparticles at the interface. When the nanoparticles are not uniformly dispersed in the PP phase, the migration itself or successful migration is greatly suppressed irrespective of the phase size. In this light, the RGT is regarded as an effective methodology for controlling the distribution of nanoparticles based on their uniform dispersion.



Fig. 2.13. TEM micrographs of PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP nanocomposites with different kinds of POE. a) POE1, b) POE2, c) POE3, d): POE4. Note that the scale bar is applied to all the samples.

# 2.3.3. Thermal conductivity

Fig. 2.14 compares the thermal conductivity of various samples prepared in this chapter, that include neat PP, PP/POE blends, and PP/Al<sub>2</sub>O<sub>3</sub>-RGT, PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT and PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP nanocomposites. The thermal conductivity of neat PP was 0.21 W/m K. It increased to 0.27 and 0.30 W/m K when 2.5 and 5.5 vol% of Al<sub>2</sub>O<sub>3</sub> nanoparticles were uniformly dispersed in the PP matrix, respectively (PP/Al<sub>2</sub>O<sub>3</sub>-RGT, see Fig. 2.3). In the case of uniform dispersion, thermal transport happens from one particle to the closest particle by crossing the organic pathway with poor thermal conductivity [53]. Accordingly, the uniform dispersion is more efficient than agglomeration [24,54–57], but not suitable to form a thermal conductive network with direct contact among nanoparticles for omitting high interfacial thermal resistance of the organic pathway [58,59].



Fig. 2.14. Thermal conductivity of various samples prepared in this chapter. The values and error bars represent the average and standard deviation of ten independent measurements per sample, respectively. Different series of samples were classified by colors, where the numbers in the x-axis correspond to the POE type except for PP/Al<sub>2</sub>O<sub>3</sub>-RGT, where the volume percentage of Al<sub>2</sub>O<sub>3</sub> is written (2.5 and 5.5 vol%). Note that the volume percentage of Al<sub>2</sub>O<sub>3</sub> was kept at 2.8 vol% for all the PP/POE/Al<sub>2</sub>O<sub>3</sub> samples.

After adding POE, the thermal conductivity was improved even when preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles were used. When compared at a similar Al<sub>2</sub>O<sub>3</sub> loading, the thermal conductivity of PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP (2.8 vol%) was 0.30– 0.33 W/m K, i.e. around 11–22% higher than that of PP/Al<sub>2</sub>O<sub>3</sub>-RGT (2.5 vol%). Considering that the dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles was superior for PP/Al<sub>2</sub>O<sub>3</sub>-RGT, different factors must account for this improvement. One

clear reason was the localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the PP phase, which in turn increased the concentration of the nanoparticles in this phase. Provided that all the Al<sub>2</sub>O<sub>3</sub> nanoparticles were localized in the PP phase, the loading to PP was assumed to be 5.5 vol%. Accordingly, PP/Al<sub>2</sub>O<sub>3</sub>-RGT (5.5 vol%) was prepared and its thermal conductivity, 0.30 W/m K, was found to be much closer to that of PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP.

The thermal conductivity was further improved when the RGT and the immiscible blend were combined to prepare PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites. Among them, PP/POE1/Al<sub>2</sub>O<sub>3</sub>-RGT achieved the highest thermal conductivity. The uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles and the small phase domain size guaranteed the successful migration of Al<sub>2</sub>O<sub>3</sub> nanoparticles to the interface and formed well-developed thermal conductive networks.

To the end, the effect of the chemical composition of POE is discussed. The thermal conductivity became higher for POE having a greater octene content (i.e. POE1,4), while a more developed network was formed for POE having higher viscosity (i.e. POE1,2). This discrepancy was understood by breaking up the thermal conductivity in the specific heat capacity and the thermal diffusivity. A higher octene content increases the specific heat capacity due to the mobility of side chains [60] (Table 2.1), while higher viscosity enhances

the thermal diffusivity (Fig. 2.15), both of which can positively contribute to the thermal conductivity.



Fig. 2.15. Thermal diffusivity of various samples preapred in this study. It can be seen that PP/POE1,2/Al<sub>2</sub>O<sub>3</sub>-RGT improved the thermal diffusivity most due to the effective migration of Al<sub>2</sub>O<sub>3</sub> nanoparticles to the interface.

# 2.4. Conclusions

In this chapter, an immiscible polymer blend strategy was employed in combination with the reactor granule technology (RGT) in order to control the selective localization of  $Al_2O_3$  nanoparticles at the interface.  $Al_2O_3$ nanoparticles thermodynamically prefer the localization at the interface between polypropylene (PP) and polyolefin elastomer (POE), but their successful migration to the interface largely depended on the initial dispersion of the nanoparticles and the phase domain size of the blends. For instance, preformed when nanoparticles were directly melt-mixed, severe agglomeration hindered the migration to the interface, while even with uniform initial dispersion, the coarse phase morphology tended to promote the agglomeration over the migration. The combination of the RGT and POE having high viscosity satisfied the prerequisites of uniform initial dispersion and a small phase size for the short migration distance, thus leading to the selective localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles at the interface, creating thermal conductive networks along the co-continuous interface, and eventually giving high thermal conductivity. The results obtained in this chapter firstly demonstrated that the distribution of nanoparticles can be controlled even for the combination of hydrophobic and hydrophilic pairs like PP and Al<sub>2</sub>O<sub>3</sub>, in which initially uniform dispersion of the nanoparticles in one phase and a small phase domain size are essential to drive successful migration of the nanoparticles. In conclusion, I successfully proposed a facile way to control the localization of filled nanoparticles at the interface in nanocomposites, and key factors thereof. The proposal is believed to be widely applicable for combination of other fillers and immiscible blends, and relevant to practical applications as electrical and thermal conductive materials.

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# Chapter 3

# Cooperative Influences of Nanoparticle Localization and Phase Coarsening on Thermal Conductivity of Polypropylene/Polyolefin Elastomer Blends

# Abstract

Polypropylene (PP)/polyolefin elastomer (POE)/aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) nanocomposites were prepared based on a reactor granule technology (RGT). Cooperation of nanoparticle migration and phase morphology evolution on the formation of thermal conductive networks was systematically studied. The utilization of RGT and POE with high viscosity enabled successful migration of insitu generated Al<sub>2</sub>O<sub>3</sub> nanoparticles at the interface of PP and POE, as combined results of uniform dispersion of Al<sub>2</sub>O<sub>3</sub> and a small phase domain size of PP. The localization of the nanoparticles at the interface also contributed to suppression of the phase coarsening during annealing. The nanoparticle migration to the interface and the retarded phase coarsening synergistically contributed to the enhanced thermal conductivity.

**Keywords:** Thermal conductivity; Nanocomposites; Immiscible blends; Nanoparticle localization; Phase coarsening

# 3.1. Introduction

With the continuous miniaturization and lightweight of electronic components, thermal conductive polymer composites have been widely studied and used in recent years due to their excellent properties such as low density, good processing ability, and resistance to corrosion [1–7]. Electrical insulating fillers with high thermal conductivity such as aluminum nitride (AIN) [8], boron nitride [9,10], silicon carbide [11], and aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) [12] are commonly introduced into polymer matrices to fabricate thermal conductive composites. Many methods have been reported to improve the thermal conductivity of the composites at a reduced loading and it has been found that the thermal conductivity greatly depends on the construction of interconnected thermal conductive networks [13–15]. In particular, the utilization of a co-continuous structure of immiscible blends is regarded as an effective way to construct such networks by the selective localization of fillers in one phase or at the interface. This strategy has been widely applied to reduce the percolation threshold in electrical conductive composites [16–18], and its application to thermal conductive composites started to appear in recent years [19,20].

The distribution of fillers in polymer blends is manipulated by thermodynamic and kinetic factors [21]. Fillers are preferentially distributed in a phase with the surface energy closer to that of fillers [22,23], where surface modification of fillers is one effective way for controlling the localization [24]. When the nanoparticles are initially distributed in the thermodynamically less favorable phase, they tend to migrate to another phase or the interface to reach the equilibrium. For instance, Gong et al. reported the migration of multi-walled carbon nanotubes (MWCNTs) in polyoxyethylene (PEO)/ethylene- $\alpha$ -octene random copolymer composites during melt compounding [25]. They confirmed that MWCNTs almost completely migrated from the PEO phase to the thermodynamically favorable copolymer phase. The migration is also influenced by kinetic factors such as the filler size, viscosity of the components, mixing procedure, shear strength, and etc. [21]. For instance, graphene nanoplates (GNPs) were kinetically trapped at the interface of polylactic acid/polystyrene (PS) immiscible blends via control of melt-compounding sequences, mixing times and shear rates, although GNPs thermodynamically prefer the PS phase [26]. Shear flow-induced hierarchical self-assembly of fillers is a prominent example of successful manipulation of the distribution and organization of fillers with the aid of kinetic control [27–29]. In Chapter 2, I found that the distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles was not only determined by the thermodynamic and kinetic factors but also greatly influenced by their initial dispersion state [7]. The uniform dispersion of nanoparticles is essential for controlling the migration of nanoparticles, while the formation of aggregates hinders the migration.

Another important parameter for the filler migration is the evolution of phase morphology in immiscible blends. The as-melt-mixed phase morphology is not in an equilibrium state and it evolves to lower interfacial free energy once the external shear is stopped. The most direct consequence of decreasing the interfacial area is the phase coalescence, leading to the increase of the phase domain size [30]. Cooperation between the filler migration and phase morphology evolution has been reported in several papers. Sun et al. studied the relationship between the phase morphology and electrical conductivity of polypropylene (PP)/PS/carbon black (CB) composites with a co-continuous structure [31]. They found that the electrical conductivity of the composites greatly increased after thermal annealing because the evolution of the co-continuous morphology drove the CB nanoparticles to selfassemble into complete nanoparticle networks. Enhancement of electrical properties by thermal annealing was also reported for a PS/poly(methyl methacrylate) (PMMA)/MWCNTs ternary system [32]. Such enhancement was caused by the coalescence of small phases into co-continuous phases and the enrichment of MWCNTs at the interface of PS/PMMA blends during the thermal annealing. Pan et al. reported an opposite phenomenon that the phase coarsening upon annealing caused the destruction of the conductive pathway of CB nanoparticles [33]. The above-introduced literatures all describe the influence of the phase morphology evolution on the filler migration, while the influence in the opposite direction, i.e. from the filler migration to the phase morphology evolution, is also known. For instance, Li et al. reported that the selective localization of nanoparticles (MWCNTs and SiO<sub>2</sub>) at the interface of polymer blends suppressed the coalescence of polymer domain and resulted in a smaller phase domain size [34,35]. Likewise, the filler migration and localization in a co-continuous structure are greatly affected by the evolution of the phase morphology, and this rationalizes a detailed investigation on

the relationship among the filler migration, phase morphology evolution, and formation of thermal conductive networks for designing thermal conductive composites.

In this chapter cooperation of the filler migration and phase morphology evolution was investigated in relation to the formation of thermal conductive networks in PP/POE/Al<sub>2</sub>O<sub>3</sub> composites with a co-continuous structure. First, the reactor granule technology (RGT) [7,36–39] was employed to facilitate the uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the PP matrix. Then, POE was blended to fabricate PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites with a co-continuous structure. The obtained nanocomposites were treated under quiescent melting annealing. The evolution of the nanoparticle localization, phase coarsening, and thermal conductivity during this process was studied comprehensively. Such study is expected to provide helpful guidance for designing conductive composite materials and engineering the phase morphology of immiscible polymer blends.

# 3.2. Experimental section

#### 3.2.1. Materials

Polypropylene reactor granule with the weight average molecular weight of 260,000, median size of 637  $\mu$ m, and pore volume of 0.56 mL/g was synthesized by catalyzed propylene polymerization. Aluminium isopropoxide (Al(O*i*Pr)<sub>3</sub>, supplied

by Sigma-Aldrich), *n*-octdecyl-3-(3',5'-di-*t*-butyl-4'-hydroxyphenyl)propionate (AO-50, donated by ADEKA), and bis(1,2,2,6,6-pentamethyl-4-piperidyl) sebacate (LA-77, donated by ADEKA) were used as a precursor, anti-oxidant and catalyst, respectively [40]. Polyolefin elastomer with the octene content of 38 wt% and the melt flow index of 1 g/10 min (measured at 2.16 kg and 190 °C based on ASTM D1238) was supplied by Dow Chemical Co. (Midland, MI, USA). Preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles with the average diameter of 15 nm (Alu, Nippon Aerosil) were used to prepare reference samples.

#### 3.2.2. Sample preparation

Nanocomposites were obtained by a masterbatch method, as schematically depicted in Fig. 3-1. First, a PP/Al<sub>2</sub>O<sub>3</sub> masterbatch was prepared as follows: 10 g of PP reactor granule was impregnated with a specific amount of Al(O*i*Pr)<sub>3</sub> dissolved in 20 mL of toluene at 50 °C under N<sub>2</sub> for 12 h. The amount of the precursor was adjusted so as to obtain different Al<sub>2</sub>O<sub>3</sub> loadings in the resultant masterbatch under the assumption of full conversion of the metal alkoxide to the oxide. After the impregnation and solvent removal in vacuo, the impregnated granule was treated at 50 °C and 100% relative humidity for 24 h to promote the hydrolysis and solidification of Al(O*i*Pr)<sub>3</sub> in the pores of PP. The obtained granule was melt-mixed at 190 °C and 100 rpm for 10 min (Micro Compounder IM5, Xplore), then a specified amount of POE was added, and blended for another 5 min. The ratio of PP

and POE was kept at 50/50 wt/wt to maintain the co-continuous phase morphology. After mixing, the product was rapidly quenched in cold water to freeze the phase morphology. The obtained samples were denoted as PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT-*x*, where *x* represents the volume fraction of Al<sub>2</sub>O<sub>3</sub> nanoparticles. Preformed Al<sub>2</sub>O<sub>3</sub> nanoparticles were also melt-mixed with PP and then with POE. These reference samples were termed as PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP-*x*.



Fig. 3.1. Schematic image of sample preparation.

Nanocomposites were compression-molded into 200  $\mu$ -thick films. A melt-mixed sample was preheated at 190 °C for 1 min, then compressed under 10 MPa for 30 s. After releasing the pressure, the sample was kept in the mold and annealed at the same temperature for different durations.

# 3.2.3. Characterization

To confirm the conversion of the precursor into Al<sub>2</sub>O<sub>3</sub>, attenuated total reflection infrared spectra (ATR-IR) (FT-IR, 100, Perkin Elmer) were acquired with a resolution of 4 cm<sup>-1</sup> and 24 scans. The Al<sub>2</sub>O<sub>3</sub> loading was determined based on the thermogravimetric analysis (TGA, Thermo plus evo, Rigaku). A specified amount of a sample was loaded into an alumina pan, followed by heating up to 600 °C under dry air flow. The residual inorganic content at 600 °C was regarded as the Al<sub>2</sub>O<sub>3</sub> content.

The dispersion, distribution and migration of Al<sub>2</sub>O<sub>3</sub> nanoparticles during the annealing process were studied by transmission electron microscopy (TEM, Hitachi H-7100, 100 kV). TEM specimens with the thickness of 100 nm were prepared by an ultramicrotome (Reichent Ultracut FCS, Reica) equipped with a diamond knife (Diatome). The phase morphology was observed using a scanning electron microscope (SEM, Hitachi S-4100) at 20 kV. A sample was cryo-fractured in liquid N<sub>2</sub>, then etched by heptane at 80 °C for 24 h to remove POE. The sample was sputtered with Pt/Pd prior to the measurements. The phase domain size was obtained by measuring the widths of rod-like holes and the diameters of circular holes in cross-sectioned SEM images as is often done in the field [31,41–44]. With the aid of Image J software, more than 100 data points were obtained to calculate the average phase domain size for each sample. Elemental mapping was also performed

using an energy dispersive X-ray analyzer (EDX, Hitachi TM 3030Plus).

The thermal diffusivity ( $\alpha$ ) of nanocomposites was determined by a temperature wave analyzer (ai-Phase mobile 1u/2, Hitachi High-Tech Science). A film sample was placed between the heater and the sensor plates. The phase delay of the temperature wave was acquired at eight frequency points within the range of 0.2–2 Hz. The thermal conductivity ( $\lambda$ ) was determined by

$$\lambda = \alpha C_p \rho \tag{3-1}.$$

The specific heat capacity ( $C_p$ ) was derived from differential scanning calorimetry (DSC, Mettler Toledo DSC-822) in the temperature range of 0–50 °C. The density ( $\rho$ ) was measured using an electronic density meter (SANSYO DME-220).

## 3.3. Results and discussion

#### 3.3.1. Distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles

When nanoparticles are introduced into immiscible polymer blends, they are usually distributed in one phase or at the interface rather than homogeneously distributed in the whole polymer blends. In an equilibrium state, the distribution of fillers is dictated by the minimum interfacial energy [45,46]. The wetting coefficient ( $\omega_a$ ) proposed by Sumita et al. [47] is often used to predict the preferential distribution of fillers in immiscible polymer blends,

$$\omega_a = \frac{\gamma_{CA} - \gamma_{CB}}{\gamma_{AB}} \tag{3-2},$$

where  $\gamma_{CA}$ ,  $\gamma_{CB}$  and  $\gamma_{AB}$  are the interfacial energies among polymer A, polymer B, and filler C. The values of  $\omega_a$  above 1, below –1, and in between mean that the fillers would preferentially be distributed in polymer B, in polymer A and at the interface, respectively. The interfacial energies  $\gamma_{ij}$  can be calculated based on the harmonic mean equation [48],

$$\gamma_{12} = \gamma_1 + \gamma_2 - 4 \left( \frac{\gamma_1^d \gamma_2^d}{\gamma_1^d + \gamma_2^d} + \frac{\gamma_1^p \gamma_2^p}{\gamma_1^p + \gamma_2^p} \right)$$
(3-3),

where  $\gamma_i$ , the surface free energy of component *i*, is divided into the dispersion and polar parts,  $\gamma_i^d$  and  $\gamma_i^p$ . These constants are listed for PP, POE, and Al<sub>2</sub>O<sub>3</sub> in Table 3.1 [49,50]. Based on these values and Eq. (3-2), the wetting coefficient was derived as -0.85, i.e. Al<sub>2</sub>O<sub>3</sub> nanoparticles were supposed to migrate to the interface, starting from the PP phase.

Materials	Total surface energy	Dispersion part	Polar part
	$\gamma (mJ/m^2)$	$\gamma^d (\mathrm{mJ/m^2})$	$\gamma^p (\mathrm{mJ/m^2})$
PP <sup>a</sup>	21.71	21.23	0.48
POE <sup>a</sup>	17.51	15.61	1.90
$Al_2O_3{}^b$	114.00	43.70	70.30

Table 3.1. Surface free energies of the filler and polymers.

<sup>a</sup> According to Ma et al. [49].

<sup>b</sup> According to Holysz et al. [50].

In order to confirm the distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles, the morphology of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites was studied by TEM, SEM, and EDX (Fig. 3.2). When the RGT was employed, Al<sub>2</sub>O<sub>3</sub> nanoparticles were selectively distributed and uniformly dispersed in the PP phase of a co-continuous structure (Fig. 3.2a). The formation of Al<sub>2</sub>O<sub>3</sub> was also confirmed by TGA and ATR-IR (Table 3.2 and Fig. 3.3). The uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles (Fig. 3.4a,b) in the PP matrix was attributed to the confinement of the precursor in the pores of reactor granule and its spontaneous conversion to the oxide. After annealing at 190 °C for 60 min, most of Al<sub>2</sub>O<sub>3</sub> nanoparticles were migrated to the interface (Fig. 3.2b). These results suggested that Al<sub>2</sub>O<sub>3</sub> nanoparticles were kinetically retained in the PP phase before annealing, and the annealing facilitated their migration to the interface in accordance with the estimation from the wetting coefficient. When the preformed nanoparticles were used, huge agglomerates were formed and the migration to the interface became much less efficient (Fig. 3.2c,d). Considering that the size of the nanoparticles was not largely different between the two protocols, the particle clusters and agglomerates (Fig. 3.4c,d) significantly hindered the migration.



Fig. 3.2. TEM (top), SEM (middle), and EDX mapping for Al (bottom) of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites: RGT-1.3 vol% a) before and b) after annealing at 190 °C for 60 min, and NP-1.3 vol% c) before and d) after annealing at 190 °C for 60 min.

Commis	Theoretical loading <sup>a</sup>	Actual loading <sup>b</sup>
Sample	(wt %)	(wt %)
PP/Al <sub>2</sub> O <sub>3</sub> -RGT-1.2 vol%	5	4.7
PP/Al <sub>2</sub> O <sub>3</sub> -RGT-2.5 vol%	10	9.5
PP/Al <sub>2</sub> O <sub>3</sub> -RGT-5.5 vol%	20	20.1
PP/Al <sub>2</sub> O <sub>3</sub> -RGT-9.1 vol%	30	29.6

Table 3.2. TGA results for PP/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites prepared by RGT.

<sup>a</sup> Theoretical loading was estimated on the assumption of the full conversion to

 $Al_2O_3.$ 

<sup>b</sup> Measured by TGA.



Fig. 3.3. ATR-IR spectra of PP/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites with different Al<sub>2</sub>O<sub>3</sub> loadings.



Fig. 3.4. TEM micrographs of PP/Al<sub>2</sub>O<sub>3</sub> masterbatch nanocomposites with different Al<sub>2</sub>O<sub>3</sub> loadings: a) RGT-2.5 vol%, b) RGT-5.5 vol%, c) NP-2.5 vol%, and d,) NP-5.5 vol%. Note that the scale bar is applied to all the images. The insets show the particles size distribution of Al<sub>2</sub>O<sub>3</sub>.

# 3.3.2. Migration of Al<sub>2</sub>O<sub>3</sub> nanoparticles and phase morphology evolution

The distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles changed after annealing and the selective localization at the interface was achieved by RGT. In order for deeper understanding, the migration of nanoparticles during the annealing process was tracked. Fig. 3.5

summarizes TEM images of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites that were annealed for different durations. When the RGT was employed, the selective localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles was eventually achieved irrespective of the loading. However, the time required for the localization at the interface was found to be shorter at a higher loading. At 0.6 vol%, a part of nanoparticles already reached the interface at 5 min, and successive localization eventually formed an Al<sub>2</sub>O<sub>3</sub> layer at 30 min. At 2.8 vol%, many Al<sub>2</sub>O<sub>3</sub> nanoparticles were found at the interface even before annealing and a continuous Al<sub>2</sub>O<sub>3</sub> layer was formed at 10 min. For the preformed nanoparticles, the migration to the interface happened, but its efficacy was quite poor. At 0.6 vol%, a few nanoparticles arrived at the interface after annealing for 10 min. With the increase of the annealing duration, more nanoparticles arrived at the interface but a continuous Al<sub>2</sub>O<sub>3</sub> layer was not eventually formed. Even at an elevated filler loading (2.8 vol%), many of Al<sub>2</sub>O<sub>3</sub> nanoparticles stayed in the PP phase after annealing for 60 min. It seems that nanoparticles hardly migrated once they were agglomerated or clustered, and their existence also hindered the migration of other particles, leading to the low efficacy. The same nanoparticle migration process was also observed by SEM images (Fig. 3.6).

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Fig. 3.5. TEM micrographs of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites with different Al<sub>2</sub>O<sub>3</sub> loadings. The filler migration was tracked along the annealing at 190 °C. Note that the scale bar is applied to the whole micrographs.


Fig. 3.6. Magnified SEM images of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites tracked along the annealing at 190 °C.

Prerequisites for achieving the selective localization are i) thermodynamic preference of the localization for given combination, ii) uniform dispersion without particle agglomeration that enables the migration of individual nanoparticles, and iii) diffusion of nanoparticles faster than phase coarsening. The first two prerequisites are satisfied in PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites, while the last prerequisite depends on the balance between the size of nanoparticles and the viscosity of the matrix. In general, the Brownian motion of a spherical nanoparticle with a radius *r* in a fluid with the viscosity  $\eta$  at temperature *T* is described by the Stokes-Einstein equation [51],

$$D_0 = \frac{k_B T}{6\pi\eta r} \tag{3-4},$$

where  $D_0$  and  $k_B$  correspond to the diffusion coefficient and the Boltzmann constant, respectively. On the one hand, the addition of nanoparticles increases the viscosity of a molten polymer [52], which can be expressed by Mooney's equation [53,54],

$$\eta_c = \eta_r exp\left(\frac{2.5f}{I-af}\right) \tag{3-5},$$

where  $\eta_c$  is the viscosity of the filled polymer,  $\eta_r$  is the viscosity of the unfilled polymer, *f* is the volume fraction of fillers, and *a* is called self-crowding factor that is related to the critical filler fraction for a steep increment in the viscosity. The size of Al<sub>2</sub>O<sub>3</sub> nanoparticles was found to be insensitive to the loading (see Fig. 3.4). According to Eq. (3-4) and Eq. (3-5), the increase of the viscosity at a higher loading would slow down the diffusion of nanoparticles. However, the observation in Fig. 3.5 dictated the opposite trend. This contradiction can be understood only when the presence of nanoparticles strongly affects the phase morphology and its evolution.

Accordingly, the evolution of the phase morphology was studied based on SEM micrographs of etched samples (Fig. 3.7). The corresponding phase domain sizes are given in Fig. 3.8. For pure PP/POE blends, the phase domain size increased from 4.1  $\mu$ m to 36.5  $\mu$ m after annealing at 190 °C for 20 min. The addition of nanoparticles caused the reduction of the phase domain size before annealing, and suppressed the phase coarsening during the annealing. The extents of the reduction

and suppression were highly dependent on the employed protocol. When preformed nanoparticles were melt-mixed at 0.6 vol%, the initial phase size became 3.5  $\mu$ m and it was enlarged to 32.6  $\mu$ m after 30 min. The large phase domain size made it difficult for nanoparticles to reach the interface. This is another reason why the migration of nanoparticles was less effective in PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP. The utilization of the RGT at the same loading made much greater impacts: 1.1  $\mu$ m at 0 min and 9.8  $\mu$ m at 30 min. The phase coarsening rate (*k*) in an immiscible polymer blend with the interfacial tension  $\sigma$  and viscosity  $\eta_c$  can be expressed by Eq. (5) [44,55],

$$k = c \frac{\sigma}{\eta_c} \tag{3-6},$$

where *c* is a dimensionless constant. According to Eq. (3-5) and Eq. (3-6), the deceleration of the phase coarsening was a result of the viscosity increment by the addition of nanoparticles, where the RGT gave larger impacts due to the uniform dispersion. When the loading was increased for the preformed nanoparticles, the phase coarsening became gradually slower, eventually 6.4  $\mu$ m after 30 min at the loading of 4.7 vol%. In the RGT case, abrupt reduction in the initial phase domain size and significant suppression of the phase coarsening were observed over 2.8 vol% (0.8 to 2.3  $\mu$ m for 60 min at 2.8 vol% and 0.7 to 1.7  $\mu$ m for 60 min at 4.7 vol%). Considering that the nanoparticle concentration was proportional to the loading, such abrupt change was plausibly ascribed by the following reasons: The viscosity of the PP phase was increased in a non-linear fashion due to the formation of

physical network structures [40], which in turn decreased the initial phase domain size and slowed down the phase coarsening. Due to a smaller phase domain size and a higher nanoparticle concentration, many Al<sub>2</sub>O<sub>3</sub> nanoparticles successfully reached the interface in a very early timing as we have observed in Fig. 3.5, and the formation of a rigid Al<sub>2</sub>O<sub>3</sub> layer at the interface plausibly prevented the phase coarsening as like a barrier layer. It is important to say that the localization of nanoparticles and phase morphology evolution strongly affect with each other. Thus, the utilization of RGT at a sufficiently high loading afforded nanocomposites bearing selectively localized and uniformly distributed thermal conductive networks.



Fig. 3.7. Evolution of the phase morphology of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites tracked along the annealing at 190 °C. Note that the scale bar is applied to all micrographs.



Fig. 3.8. Phase domain size of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites as a function of a) annealing time and b) Al<sub>2</sub>O<sub>3</sub> loading.

### 3.3.3 Thermal conductivity

Cooperative influences of nanoparticle localization and phase morphology evolution on the thermal conductivity of the nanocomposites were studied. Fig. 3.9 summarizes the thermal conductivity of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites prepared by different protocols as a function of annealing time. At a fixed Al<sub>2</sub>O<sub>3</sub> loading, PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT nanocomposites showed obviously better thermal conductivity than PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP nanocomposites. Besides, the dependence of the thermal conductivity on the annealing time was largely different between the two series and among different loadings. In the case of PP/POE/Al<sub>2</sub>O<sub>3</sub>-RGT at 0.6 and 1.3 vol%, the thermal conductivity increased along annealing time in the first 30 min and decreased slightly with the further increase of the annealing time. When the loading was increased over 2.8 vol%, the thermal conductivity increased with the increase of annealing time and it never decreased by elongated annealing. In the case of PP/POE/Al<sub>2</sub>O<sub>3</sub>-NP nanocomposites, the thermal conductivity once reached the maximum at 10-30 min and then slowly decreased by the further annealing.



Fig. 3.9. Thermal conductivity of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites annealed at 190 °C for different durations.

As stated earlier, the annealing can cause both positive and negative influences on the thermal conductivity. When preformed nanoparticles were employed or when the loading was not sufficiently high for RGT, the phase coarsening cannot be sufficiently retarded (Fig. 3.8). In the early stage of the annealing, the migration of Al<sub>2</sub>O<sub>3</sub> nanoparticles to the interface led to enhancement of the thermal conductivity due to the formation of thermal conductive pathways. The extent of the increment was greater for RGT as most of the nanoparticles were successfully migrated due to the initially uniform dispersion. When the annealing was extended, the negative effect of the phase coarsening was not anymore compensated by further localization. On the other hand, when RGT was employed at a high loading, the phase coarsening was almost completely suppressed. As a result, the thermal conductivity of the nanocomposites monotonously increased by the annealing. Rapid evolution below 10 min corresponded to the time required to form connected networks at the interface, while the slower evolution over 10 min corresponded to the network thickening.

Based on the results and analysis above, the relationship among the evolutions of nanoparticle localization, phase morphology, and thermal conductive networks is illustrated in Fig. 3.10. The uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles modifies viscoelastic properties of the PP phase and brings about smaller phase domains (Fig. 3.10a). The distance of migration from the initial position to the interface is hence greatly shortened and migration completes in a shorter duration (Fig. 3.10a'). Thus, localized nanoparticles at the interface form thermal conductive networks. These networks also act as a barrier to suppress the phase coarsening. The particle

agglomerates and clusters disturb all the above mechanisms: Smaller viscoelastic modulation, less efficient migration, and less suppressed phase coarsening. The combination of these factors makes the thermal conductivity of resultant nanocomposites much poorer (Fig. 3.10b,b').



Fig. 3.10. Schematic representation of the network and phase morphology evolution in PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites.

## 3.4. Conclusions

The controlled migration of nanoparticles within a co-continuous structure of

immiscible polymer blends is a key issue to design thermal conductive composite materials. In this chapter, the relationship among nanoparticle migration, phase morphology evolution, and the thermal conductivity were reported based on a PP/POE/Al<sub>2</sub>O<sub>3</sub> ternary system. Al<sub>2</sub>O<sub>3</sub> nanoparticles spontaneously migrated from the PP phase to the interface, where the uniform dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the PP phase was essential. Important cooperation between nanoparticle localization and phase morphology evolution was evidenced. Nanoparticles that were uniformly dispersed in the PP phase decreased the phase domain size to facilitate successful migration to the interface. Further, the formation of nanoparticle networks decelerated the phase coarsening during annealing. By employing the RGT at a sufficiently high Al<sub>2</sub>O<sub>3</sub> loading, uniformly distributed thermal conductive networks were afforded. These findings are believed important not only in designing conductive composite materials but also in engineering the phase morphology of immiscible polymer blends for other applications.

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# Chapter 4

# **Design of Continuous Segregated Polypropylene/Al<sub>2</sub>O<sub>3</sub>**

## Nanocomposites and Impact of Controlled Al<sub>2</sub>O<sub>3</sub>

**Distribution on Thermal Conductivity** 

## Abstract

Control of nanoparticle distribution in polymer matrices is a key factor for designing highly conductive nanocomposites. Here, polypropylene (PP)/aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) nanocomposites with a continuous segregated structure were designed. Al<sub>2</sub>O<sub>3</sub> nanoparticles were initially distributed in the polyolefin elastomer (POE) phase of PP/POE/Al<sub>2</sub>O<sub>3</sub> with a co-continuous structure. Selective extraction of the POE phase provided a porous PP scaffold, whose pore walls were covered by deposited Al<sub>2</sub>O<sub>3</sub> nanoparticles. Subsequent compression molding made the porous scaffold tightly compacted to form uniform and dense thermal conductive networks. The thermal conductivity was compared among nanocomposites having three different types of Al<sub>2</sub>O<sub>3</sub> distribution. It was found that the continuous segregated distribution was far the most effective for improving the thermal conductivity, where 1.07 W/m K was achieved at an Al<sub>2</sub>O<sub>3</sub> loading of 27.5 vol%.

**Keywords**: Thermal conductivity; Continuous segregated network; Nanocomposite; Aluminum oxide; Polypropylene

## 4.1. Introduction

In the past decade, thermal conductive polymer composites have attracted great attention due to the continuous integration and miniaturization of microelectronic devices [1–3]. The thermal conductivity of neat polymer is often insufficient, and

thermally conductive and insulating fillers such as boron nitride (BN) [4], aluminum nitride [5], silicon carbide [6], and aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) [7] are usually introduced to improve the thermal conductivity and simultaneously meet the stringent requirements for electrical insulation. However, conventional methods such as melt mixing, solution mixing, or powder mixing [8] usually lead to random distribution of fillers in the polymer matrix. Therefore, desired thermal conductivity can be obtained only at an unacceptably high filler loading, which causes problems in density, mechanical properties, processability, and fabrication cost [9,10].

It is recognized that the thermal conductivity is determined not only by the choice of fillers, but also interfacial interaction in relation to phonon scattering [11]. The inclusion of fillers of a finite size necessarily generates interfaces. This discontinuity or interfacial mismatch causes phonon scattering [12], i.e. even with the same type and amount of filler, the resultant thermal conductivity of composites largely depends on the design of interfaces. For example, when filler particles are randomly distributed in a polymer matrix, the interfacial connection between the polymer and filler is important. Surface modification of fillers is often adopted to improve the interfacial connection between the polymer and fillers. Shen et al. reported that modification of BN by polydopamine greatly improved the thermal conductivity of polyvinyl alcohol (PVA) to 5.4 W/m K at 10 vol%, which was 63% higher than that of PVA/unmodified BN [13]. Here, it must be noted that surface modification does not necessarily improve the thermal conductivity. Gulotty et al. studied the effect of the surface modification of single-wall carbon nanotubes on the thermal conductivity of polymer nanocomposites [14]. It was found that the intrinsic thermal conductivity of fillers decreases with the modification, and this negative effect sometimes overcomes the positive effect of improved interfacial connection [15].

A more appropriate way to obtain high thermal conductivity is to form a continuous network through selective distribution of fillers [16-18]. Unlike the propagation of electrons in electric conductive composites, where they can "hop" between fillers through tunneling [19,20], a direct contact between fillers is required for thermal conduction. Many strategies have been proposed to construct continuous thermal conductive networks. In one strategy, a co-continuous morphology of immiscible polymer blends is exploited to control the selective localization of fillers. A volume exclusion effect increases the local concentration of fillers and improves the connectivity of the networks [20–22]. In particular, the selective localization at the interface of a co-continuous structure is one of the most effective ways to form thermal conductive networks at a lower filler loading [23]. However, this strategy relies on a subtle balance of surface energies among polymers and fillers [9]. Moreover, the distribution of fillers is also affected by kinetic factors such as the viscosity, the mixing procedure, and the filler size. In the previous two chapters, I applied a novel reactor granule technology [7,24,25] for the fabrication of thermally

conductive polypropylene (PP)-based nanocomposites. I found that successful migration of nanoparticles to the interface is achieved only when nanoparticles are initially uniformly dispersed, and their migration is faster than the phase coarsening [26,27]. In recent literature, an alternative strategy of forming continuous networks based on a segregated structure is presented for electrically conductive polymer composites [28,29]. Typically, micron-sized polymer particles are physically mixed with fillers in a solid state, then the mixture is hot pressed. This restricts the placement of fillers at the interfaces of particles, thereby forming continuous and dense networks to greatly lower the percolation threshold. A potential concern of this strategy is imperfection of networks, e.g. physical blending may lead to non-uniform coating of polymer particles by fillers, random packing of particles may remain voids even after the hot-pressing, and so on [30–32].

This chapter attempts to enhance the thermal conductivity by constructing a continuous segregated structure in PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites. Al<sub>2</sub>O<sub>3</sub> nanoparticles were adopted due to their abundance and electrically insulating nature. Moreover, this spherical thermal conductive filler is sensitive to the uniformity and the density of the networks so that it can reflect the efficiency of the proposed strategy. Firstly, Al<sub>2</sub>O<sub>3</sub> nanoparticles were selectively distributed in the POE phase of PP/POE blends with a co-continuous structure. After the selective extraction of the POE phase, a porous PP scaffold with continuous channels was formed, where Al<sub>2</sub>O<sub>3</sub>

nanoparticles were uniformly deposited on the pore walls. The obtained porous scaffold was compression-molded to form a continuous segregated structure. In this unique structure, continuous and dense thermal conductive networks were formed and uniformly distributed in the PP matrix, thus achieving high thermal conductivity at a low loading of Al<sub>2</sub>O<sub>3</sub> nanoparticles.

### 4.2. Experimental section

#### 4.2.1. Materials

Isotactic polypropylene as a matrix ( $M_w = 2.6 \times 10^5$  g/mol, mmm = 98 mol%) was synthesized by propylene polymerization using a Mg(OEt)<sub>2</sub>-based Ziegler-Natta catalyst. Polyolefin elastomer with a 1-octene content of 38 % and a moony viscosity of 23 MU (ML 1+4 at 121 °C acquired based on ASTM D1646) was selected. Al<sub>2</sub>O<sub>3</sub> nanoparticles with the average diameter of 15 nm were purchased from Nippon Aerosil Co. Ltd.

#### 4.2.2. Sample preparation

The fabrication procedure of a continuous segregated structure is summarized in Fig. 4.1. First, Al<sub>2</sub>O<sub>3</sub> nanoparticles were melt-mixed with POE at 190 °C and 100 rpm for 10 min (Micro compounder IM5, Xplore). Then, PP was added and mixed for another 10 min under the same condition. The mixture was rapidly quenched in cold water kept at 0 °C to freeze the phase morphology. The product was compression-molded at 190 °C and 10 MPa for 5 min to obtain films with a cocontinuous structure, where Al<sub>2</sub>O<sub>3</sub> nanoparticles were aimed to selectively distribute in the POE phase (Fig. 4.1a). The obtained samples are denoted as PP/POE/Al<sub>2</sub>O<sub>3</sub>. The volume fraction of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the nanocomposites was varied from 1.2 to 18.6 vol%, while the weight ratio of PP and POE was kept at 50/50 to maintain the co-continuous phase morphology. The above-prepared PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposite films were immersed in heptane at 80 °C for 24 h to extract the POE phase. The POE phase was selectively replaced by heptane to form a PP/solvent/Al<sub>2</sub>O<sub>3</sub> ternary system (Fig. 4.1b). Solvent replacement in polymer blends with a co-continuous structure has been reported in a PP/polystyrene/carbon black system [32,33]. By the solvent evaporation,  $Al_2O_3$  nanoparticles were driven to the original PP/POE interface and led to the conformal Al<sub>2</sub>O<sub>3</sub> coating (Fig. 4.1c). Thus obtained porous films were vacuum dried at 60 °C for 24 h, and subjected to the second compression molding at 160 °C and 10 MPa for 5 min to compact the pores and Al<sub>2</sub>O<sub>3</sub> layers tightly. PP was not completely melted at 160 °C so that the migration of Al<sub>2</sub>O<sub>3</sub> nanoparticles into the matrix was restricted [32,34]. Through such a method, PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with a continuous segregated structure (denoted as s-PP/Al<sub>2</sub>O<sub>3</sub>) were prepared (Fig. 4.1d).

As a reference, PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with random distribution (termed r-

PP/Al<sub>2</sub>O<sub>3</sub>) were also prepared by melt-mixing PP and Al<sub>2</sub>O<sub>3</sub> nanoparticles at 190 °C and 100 rpm for 10 min. The product was compression-molded into films at 190 °C and 10 MPa for 5 min.



Fig. 4.1. Procedure for the construction of a continuous segregated structure. a) PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites with a co-continuous structure, b) PP/heptane/Al<sub>2</sub>O<sub>3</sub> ternary system with released Al<sub>2</sub>O<sub>3</sub> nanoparticles after the selective extraction of the POE phase; c) continuous porous PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with Al<sub>2</sub>O<sub>3</sub> nanoparticles assembled on the pore walls; d) PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with a continuous segregated structure.

### 4.2.3. Characterization

The morphology of the nanocomposites was observed by a scanning electron microscope (SEM, Hitachi S-4100) with an acceleration voltage of 20 kV. Samples were cryo-fractured in liquid N<sub>2</sub> and the fractured surface was sputter-coated with platinum and palladium prior to the measurement. Elemental mapping was also performed using an energy dispersive X-ray analyzer (EDX, Hitachi 3030Plus). To observe the distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles, a transmission electron microscope (TEM, Hitachi H-7100) was used at an acceleration voltage of 100 kV. TEM specimens with the thickness of 100 nm were prepared by an ultramicrotome (Reichent Ultracut FCS, Reica) equipped with a diamond knife (Diatome).

A weight loss measurement was performed to evaluate the continuity of the POE phase. PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposite films were immersed in heptane at 80 °C for 24 h to extract the POE phase. After the extraction, the films were washed and dried at 60 °C for 24 h. The continuity of the POE phase was evaluated by comparing the weight loss by the extraction with the initial weight of POE using Eq. (4-1).

$$%continuity = \frac{(weight of POE)_{initial} - (weight of POE)_{final}}{(weight of POE)_{initial}}$$
(4-1).

The Al<sub>2</sub>O<sub>3</sub> loading in the porous PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites was measured by thermal gravimetric analysis (TGA, Thermo plus evo, Rigaku) to study the remaining amount of Al<sub>2</sub>O<sub>3</sub> nanoparticles after the extraction of the POE phase. A sample was heated from 25 to 600 °C under dry air at a heating rate of 10 °C/min. The residual inorganic content at 600 °C was regarded as the Al<sub>2</sub>O<sub>3</sub> content.

The thermal diffusivity ( $\alpha$ ) of nanocomposites was measured by a temperature wave analyzer (ai-Phase mobile 1u/2, Hitachi High-Tech Science). A sample was sandwiched between the heater and the sensor plates. The phase delay in the temperature wave was measured at eight frequency values within a range of 0.2–2 Hz. The thermal conductivity ( $\lambda$ ) was derived from

$$\lambda = \alpha \ C_p \ \rho \tag{4-2}$$

The specific heat capacity ( $C_p$ ) at room temperature was determined using a differential scanning calorimeter (DSC, Mettler Toledo DSC-822). The density was measured by an electronic density meter (SANSYO DME-220).

Mechanical properties of nanocomposites were measured by a universal testing machine (Instron 3365) with a crosshead speed of 5 mm/min at room temperature. Dumbbell-shaped specimens were cut out from film samples.

## 4.3. Results and discussion

#### 4.3.1. Construction of continuous segregated structure

When nanoparticles are introduced into immiscible polymer blends, they tend to selectively distribute in one phase or at the interface toward a thermodynamic equilibrium. A prerequisite for the construction of a continuous segregated structure is that Al<sub>2</sub>O<sub>3</sub> nanoparticles are initially confined in the POE phase or at the interface,

while the selective distribution at the interface was estimated from the wetting coefficients [35]. Accordingly, Al<sub>2</sub>O<sub>3</sub> nanoparticles were firstly melt-mixed with POE prior to the addition of PP in order to prohibit kinetic entrapment in the PP phase. The morphology of the PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites was observed by TEM and SEM (Fig. 4.2). As shown in Fig. 4.2a,b, two phases existed in a mutually interpenetrated manner, corresponding to a typical co-continuous structure. The preservation of the co-continuous structure with the addition of Al<sub>2</sub>O<sub>3</sub> nanoparticles was verified by evaluating the continuity of the POE phase according to Eq. (4-1) (Fig. 4.3). For all the PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites, the continuity of the POE phase was maintained at over 95%, indicating that the morphology of all the samples was regarded co-continuous. Al<sub>2</sub>O<sub>3</sub> nanoparticles were distributed in one phase and at the interface, while the other phase contained virtually no nanoparticles. For further confirmation, the POE phase was selectively extracted and the crosssectional morphology of the etched sample was observed by SEM. As shown in Fig. 4.2c,d, the cross-sectioned surfaces of PP were smooth while Al<sub>2</sub>O<sub>3</sub> nanoparticles were exclusively distributed on the surface of the pores remained after etching the POE phase. This fact indicates that Al<sub>2</sub>O<sub>3</sub> nanoparticles were barely distributed in the PP phase.



Fig. 4.2. a,b) TEM micrographs of a PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposite with the Al<sub>2</sub>O<sub>3</sub> loading of 5.4 vol% before etching. c,d) SEM micrographs of the same sample after etching the POE phase. The measurements were performed at different magnifications.



Fig. 4.3. Continuity of the POE phase as a function of Al<sub>2</sub>O<sub>3</sub> loading.

The successful coating of pore walls by Al<sub>2</sub>O<sub>3</sub> nanoparticles in the continuous

porous structure is a next key step to construct the continuous segregated structure. Fig. 4.4 shows the SEM micrographs of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites after etching the POE phase. All the samples presented a porous structure and many continuous channels interpenetrated with the remaining PP phase. After the extraction of the POE phase and solvent evaporation, the pore surfaces were coated with Al<sub>2</sub>O<sub>3</sub> nanoparticles. When the Al<sub>2</sub>O<sub>3</sub> loading was low, a part of pore surfaces looked smooth, suggesting incomplete coating (Fig. 4.4a). Along with the increase in the loading, the smooth surfaces disappeared and the coating layer became denser and thicker (Fig. 4.4b–d).



Fig. 4.4. SEM micrographs of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites after the extraction of the POE phase with different Al<sub>2</sub>O<sub>3</sub> loadings: a, a') 1.2 vol%, b,b') 2.5 vol%, c,c') 5.4 vol%, and d,d') 8.9 vol%. The loading refers to the volume fraction of Al<sub>2</sub>O<sub>3</sub> nanoparticles in PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites before etching the POE phase. Note that the scale bar is applied to individual entire rows.

It was considered that a part of Al<sub>2</sub>O<sub>3</sub> nanoparticles could be extracted with heptane. In order to confirm the retention of Al<sub>2</sub>O<sub>3</sub> nanoparticles, the Al<sub>2</sub>O<sub>3</sub> loading in the remained porous PP scaffold was examined with TGA. The results are shown in Fig. 4.5. The residual weight at 600 °C was regarded as the weight of retained Al<sub>2</sub>O<sub>3</sub> nanoparticles, and it was compared with the theoretical Al<sub>2</sub>O<sub>3</sub> content that was derived under the assumption that nanoparticles were fully retained. The results are summarized in Table 4.1. The residual contents of Al<sub>2</sub>O<sub>3</sub> nanoparticles were 8.8, 16.8, 30.8, 42.8, 53.5, and 61.9 wt% for nanocomposites with the theoretical contents of 9.5, 18.2, 33.3, 45.8, 57.1 and 66.2 wt%, respectively. The retention percentage was higher than 92% for all the samples, suggesting that the solvent extraction and subsequent evaporation were effective for the deposition of Al<sub>2</sub>O<sub>3</sub> nanoparticles. It is noted that the thermal decomposition of porous PP/Al<sub>2</sub>O<sub>3</sub> was slower than that of neat PP at lower loadings while faster at higher loadings. This trend is likely attributed to competition between the thermal insulation by pores and acid-catalyzed decomposition of PP.



Fig. 4.5. TGA curves for porous  $PP/Al_2O_3$  nanocomposites after etching the POE phase. The loading in the figure refers to the volume fraction of  $Al_2O_3$  nanoparticles in the corresponding  $PP/POE/Al_2O_3$  nanocomposites before etching.

Sample <sup>a</sup>	Theoretical Al <sub>2</sub> O <sub>3</sub>	Residual weight	Retention
	content (wt%) <sup>b</sup>	at 600 °C (wt%)	percentage (%) <sup>c</sup>
0 vol%	0	0	n.a.
1.2 vol%	9.5	8.8	92.6
2.5 vol%	18.2	16.8	92.3
5.4 vol%	33.3	30.8	92.5
8.9 vol%	45.8	42.9	93.7
13.2 vol%	57.1	53.5	93.7
18.6 vol%	66.2	61.9	93.5

Table 4.1. Retention of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the pores.

<sup>a</sup> The sample names correspond to PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites with specified Al<sub>2</sub>O<sub>3</sub> loadings. Note that the TGA measurements were performed after etching the

POE phase.

<sup>b</sup> The theoretical Al<sub>2</sub>O<sub>3</sub> content was calculated under the assumption that Al<sub>2</sub>O<sub>3</sub> nanoparticles were perfectly retained in the porous PP scaffold after etching the POE phase.

 $^{\circ}$  The retention percentage refers to the ratio of the residual weight at 600  $^{\circ}$ C with respect to the theoretical Al<sub>2</sub>O<sub>3</sub> content.

The tight compaction of porous PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites is essential to enhance the interconnection of Al<sub>2</sub>O<sub>3</sub> nanoparticles and avoid defects that can cause phonon scattering. Before the compression, the appearance of the porous films was white due to the presence of pores. After the second compression, the pores were tightly compacted and the transparency of the film obviously increased (Fig. 4.6a). The situation was similar for different loadings, while the transparency of the compressed films slightly decreased along with the loading due to the thickening of the Al<sub>2</sub>O<sub>3</sub> networks (Fig. 4.7). From SEM micrographs, the obtained s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites were tightly compressed and voids were hardly observed (Fig. 4.6b-e). The formation of Al<sub>2</sub>O<sub>3</sub> networks became more and more evident along with the loading. EDX mapping was also conducted to confirm the formation of Al<sub>2</sub>O<sub>3</sub> networks (Fig. 4.8). The Al<sub>2</sub>O<sub>3</sub> networks were excluded from the PP phase to form continuous segregated networks.


Fig. 4.6. a) Optical images of porous PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites and s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites obtained after the second compression (for 4.5 vol%). b,c,d,e) SEM micrographs of s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with different Al<sub>2</sub>O<sub>3</sub> loadings: b, b') 2.2 vol%, c,c') 4.5 vol%, d,d') 9.4 vol%, and e,e') 14.9 vol%. The loading was obtained from the TGA results in Table 4.1. Note that the scale bar is applied to individual entire rows.



Fig. 4.7. Optical images of porous PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites and s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites obtained after the second compression. Note that the scale bar is applied to all images.



Fig. 4.8. EDX elemental mapping for Al in s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with different loadings: a) 2.2 vol%, b) 4.5 vol%, c) 9.4 vol%. Note that the scale bar is applied to all the samples.

The continuous segregated structure was further studied by TEM micrographs (Fig. 4.9). The results were basically consistent with those of SEM and EDX, while a few additional observations were obtained. The second compression molding shortened the spacing between Al<sub>2</sub>O<sub>3</sub> nanoparticles. Some of nanoparticles overlapped with each other and this tendency became more evident at a higher loading. On the other hand, the un-overlapped areas were necessarily filled by PP. It was considered that a fraction of PP melted at the temperature of the second compression molding, and it filled interparticle gaps of Al<sub>2</sub>O<sub>3</sub> nanoparticles that were not eliminated only by the compression. Indeed, when the compression was made at a lower temperature, a transparent film was hardly obtained. The selection of an appropriate temperature is important not only to fill the interparticle gaps but also to maintain the integrity of the formed networks.



Fig. 4.9. TEM micrographs of s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with different loading: a) 2.2 vol%, b) 4.5 vol%, c) 9.4 vol%, and d) 14.9 vol%. a')–d') are magnified SEM micrographs. Note that the scale bar is applied to individual entire rows.

#### 4.3.2. Thermal conductivity

Fig. 4.10a depicts the measured thermal conductivity for three types of nanocomposites having different Al<sub>2</sub>O<sub>3</sub> distributions: s-PP/Al<sub>2</sub>O<sub>3</sub>, where Al<sub>2</sub>O<sub>3</sub> nanoparticles formed continuous segregated networks (Fig. 4.10b); PP/POE/Al<sub>2</sub>O<sub>3</sub>, where Al<sub>2</sub>O<sub>3</sub> nanoparticles were selectively distributed in the POE phase (Fig. 4.10c); r-PP/Al<sub>2</sub>O<sub>3</sub>, where Al<sub>2</sub>O<sub>3</sub> nanoparticles were randomly distributed (Fig. 4.10d). It must be noted that the random distribution indicates a non-selective distribution, not the random placement of individual nanoparticles. As shown in Fig. 4.10a, the thermal conductivity value and its development along the Al<sub>2</sub>O<sub>3</sub> loading significantly depended on the distribution type. The random distribution (r-

PP/Al<sub>2</sub>O<sub>3</sub>) led to the lowest thermal conductivity at a given loading. In this case, Al<sub>2</sub>O<sub>3</sub> nanoparticles were mutually most separated and the phonon transport happened from one particle to the closest particle by crossing the organic pathway with poor thermal conductivity. The presence of many huge agglomerates was also regarded negative in terms of uniform thermal conduction. The selective distribution of Al<sub>2</sub>O<sub>3</sub> nanoparticles in the POE phase (and at the interface) greatly improved the thermal conductivity of the nanocomposites (PP/POE/Al<sub>2</sub>O<sub>3</sub>). This is a natural consequence of the concentration of Al<sub>2</sub>O<sub>3</sub> nanoparticles in one phase of a cocontinuous structure, which shortens the interparticle distance and raises a possibility to form overlapping nanoparticles. The best thermal conductivity was attained for the continuous segregated distribution (s-PP/Al<sub>2</sub>O<sub>3</sub>), where the connection between Al<sub>2</sub>O<sub>3</sub> nanoparticles was significantly improved as shown in Fig. 4.6. The thermal conductivity reached 1.07 W/m K at a filler loading of 27.5 vol%, while 0.5–0.8 W/m K at 36 vol% was reported for epoxy/Al<sub>2</sub>O<sub>3</sub> nanocomposites [36].



Fig. 4.10. a) Thermal conductivity of nanocomposites with different distributions of Al<sub>2</sub>O<sub>3</sub> nanoparticles and fitting of the thermal conductivity based on the Agari's model. TEM micrographs of these nanocomposites are shown in the right side: b) s-PP/Al<sub>2</sub>O<sub>3</sub> (9.4 vol%), c) PP/POE/Al<sub>2</sub>O<sub>3</sub> (8.9 vol%), and d) r-PP/Al<sub>2</sub>O<sub>3</sub> (10.5 vol%). The same scale bar is applied to all the images.

To obtain further insights into the impact of the distribution type on the thermal conductivity of nanocomposites, the experimentally obtained relationships between the thermal conductivity and the  $Al_2O_3$  loading were fit to a theoretical model proposed by Agari et al. [8,37]. This model was derived by generalizing parallel and series conduction with empirical correction factors,

$$log\lambda = V_f C_2 log\lambda_f + (1 - V_f) log(C_1 \lambda_p)$$
(4-3),

where  $\lambda$ ,  $\lambda_p$ , and  $\lambda_f$  correspond to the thermal conductivity of nanocomposites, polymer, and filler.  $V_f$  is the volume fraction of filler.  $C_I$  is a factor related to the effect of filler on the secondary structure of polymer, and the value of  $C_1 = 1$ indicated that fillers have no effect on the polymer structure.  $C_2$  is a factor for expressing the ease of the formation of conductive filler networks. The  $C_2$  value is supposed to vary between 0 and 1. When it is closer to 1, it means that conductive networks are more easily formed. Thus, the distribution of nanoparticles affects the thermal conductivity of nanocomposites through the  $C_2$  value. This model is suitable for systems filled with spherical particles at relatively high filler loadings. Fig. 4.10a shows the fitting results for three types of nanocomposites: s-PP/Al<sub>2</sub>O<sub>3</sub>, PP/POE/Al<sub>2</sub>O<sub>3</sub> and r-PP/Al<sub>2</sub>O<sub>3</sub>. The thermal conductivity of the nanocomposites was well fit to the equation irrespective of the distribution type, while the  $C_{1,2}$  values were sensitive to the distribution type. In particular, the  $C_2$  value for r-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites was 0.25, indicating the least effective network formation when randomly distributed. The  $C_2$  value increased dramatically when the distribution was controlled. In particular, the strategy of the continuous segregated structure afforded the highest  $C_2$  value of 0.84. This result clearly revealed the effectiveness of the strategy in constructing thermal conductive networks. The  $C_1$  value for PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites was the closest to 1, indicating that almost no effect was caused on matrices in this system. The  $C_1$  value for r-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites

was higher than that of PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites. In r-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites, Al<sub>2</sub>O<sub>3</sub> nanoparticles were randomly distributed in the whole area rather than selectively distributed at the interface, so that fillers could affect the polymer structure more. For s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites, the second compression-molding and the barrier effect of filler networks are believed the main reasons that affected the polymer structure.

### 4.3.3. Mechanical properties

It is known that the introduction of a segregated structure often accompanies significant deterioration of mechanical properties due to the presence of defects and weak interfacial interaction [38]. Accordingly, mechanical properties of s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites were measured and compared with those of r-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites (Fig. 4.11). In the case of r-PP/Al<sub>2</sub>O<sub>3</sub>, the particle agglomeration as well as weak interfacial interaction led to significant deterioration in the tensile strength and tensile modulus. Unlike this, the s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites exhibited enhanced tensile strength and tensile modulus. This enhancement was rather unexpected in a sense that most of previously reported segregated structures caused significant deterioration of mechanical properties. The utilization of a co-continuous template as well as the lack of fatal defects would explain the enhancement.



Fig. 4.11. Tensile properties of s-PP/Al<sub>2</sub>O<sub>3</sub> and r-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites: a) Tensile strength, b) tensile modulus, and c,d) elongation at break.

### 4.3.4. Continuous segregated structure based on reactor granule technology

In Chapters 2 and 3, it was found that the phase domain size of PP/POE blends was significantly reduced by employing the reactor granule technology (RGT), which was thought positive in terms of sample uniformity and uniform temperature distribution. Here, I attempted to apply the RGT to design the continuous structure for achieving more uniformly distributed thermal conductive networks.

First, Al<sub>2</sub>O<sub>3</sub> nanoparticles were melt-mixed with POE at 190 °C and 100 rpm for

10 min. Then, Al(O*i*Pr)<sub>3</sub>-impregnated PP granule was added and mixed for another 10 min under the same condition. The impregnated PP was prepared according to the same procedure described in the experiment part of Chapter 2, where the amount of the Al(O*i*Pr)<sub>3</sub> precursor was determined so as to obtain the Al<sub>2</sub>O<sub>3</sub> loading of 5, 10, and 20 wt% in resultant PP/Al<sub>2</sub>O<sub>3</sub> samples. The obtained mixture was compressed into films with the same compression condition for the PP/POE/Al<sub>2</sub>O<sub>3</sub> samples, and the films are denoted as PP/POE/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP). Finally, films with a continuous segregated structure were fabricated according to the same procedure. They are denoted as s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP).

The morphology of s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) was studied by TEM micrographs (Fig. 4.12a,b). It can be observed that Al<sub>2</sub>O<sub>3</sub> nanoparticles were selectively localized and form a connected network structure throughout the whole composites. The second compression molding shortened the spacing between Al<sub>2</sub>O<sub>3</sub> nanoparticles. Some of nanoparticles overlapped with each other and this tendency became more evident at a higher loading (Fig. 4.12a',b'). Compared to s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites, the number of thermal conductive paths was obviously improved in s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) (Fig. 4.12c,d).



Fig. 4.12. TEM micrographs of s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) nanocomposites with different loadings: a,a') 4.5 vol%, and b,b') 15.3 vol%. c,c') and d,d') are TEM micrographs of s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with the Al<sub>2</sub>O<sub>3</sub> loading of 4.5 vol% and 14.9 vol%.

Fig. 4.13 shows the thermal conductivity of s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) and s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites. It increased with the increase of the Al<sub>2</sub>O<sub>3</sub> loading for two kinds of nanocomposites, while s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) nanocomposites presented slightly higher thermal conductivity. This improvement is believed to be originated from the increased number of thermal conductive paths. When the RGT was applied to fabricate PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites, the phase domain size was significantly reduced as shown in Fig. 4.14. This small-sized phase domains led to the formation of more thermal conductive paths in s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) (Fig. 4.12a,b). On the other hand, the increased number of thermal conductive paths resulted in the reduced density of individual paths. Therefore, individual thermal conductive paths were thought to be denser in s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites.



Fig. 4.13. Thermal conductivity of  $s-PP/Al_2O_3-(RGT+NP)$  and  $s-PP/Al_2O_3$  nanocomposites.



Fig. 4.14. SEM micrographs of a,b,c) PP/POE/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) and d,e,f) PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites after etching the POE phase. a,d): 1.2 vol%, b,e): 2.5 vol%, c,f) 5.4 vol%). g) is the phase domain size obtained by the analysis of the TEM images.

From the thermal conductivity comparison between s-PP/Al<sub>2</sub>O<sub>3</sub>-(RGT+NP) and s-PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites, it can be known that both the number and density of thermal conductive paths are important for the improvement of thermal conductivity. At a fixed filler loading, the increase in the number of thermal conductive paths necessarily accompanies less dense paths. These two compete in determining the thermal conductivity. Even though great enhancement of thermal conductivity was not achieved by the introduction of RGT to the continuous segregated structure, the obtained results are useful for high thermal conductive nanocomposites design. The RGT provides an opportunity to control the balance between the number and density of thermal conductive paths for desired applications.

### 4. Conclusions

In this chapter, thermally conductive PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites were designed on the basis of a continuous segregated structure. By controlling the mixing order and thermodynamic preference, Al<sub>2</sub>O<sub>3</sub> nanoparticles were preferentially distributed in the POE phase of PP/POE blends with a co-continuous morphology. Solvent-aided extraction of the POE phase formed a porous PP scaffold with its pore walls coated by Al<sub>2</sub>O<sub>3</sub> nanoparticles. Subsequent compression of the porous scaffold led to the compaction of the porous structure and the densification of the nanoparticle packing, thereby affording PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with continuous segregated networks. The phase morphology and Al<sub>2</sub>O<sub>3</sub> distribution were comprehensively studied by TEM, SEM, and TGA. By comparing the thermal conductivity of nanocomposites having three different types of Al<sub>2</sub>O<sub>3</sub> distribution, it was revealed that the continuous segregated distribution gave far the most efficient improvement. Especially, the thermal conductivity of 1.07 W/m K at 27.5 vol% deserves special notice when isotropic nanoparticles are employed. Combination of the RGT with a continuous segregated structure offered an opportunity to control the balance between the number and density of thermal conductive paths. The proposed protocol is directly applicable for designing highly thermal and electric conductive materials, and would be useful for any other applications which require connected networks of fillers.

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# Chapter 5

## **General Conclusion**

Thermal conductive polymer composites have become a hot research topic along with the rapid miniaturization and high degree of integration of electronic devices. The current research on this topic has mostly assumed the use of specific thermal conductive fillers or methods to improve the thermal conductivity of polymer composites, which are generally only suitable for limited polymer matrices and very expensive. Polyolefin, as one of the most commonly used thermoplastics with balanced properties and price, is very promising to be applied in this field if its thermal conductive fillers. In this thesis, I employed polypropylene and Al<sub>2</sub>O<sub>3</sub> nanoparticle as the polymer matrix and thermal conductive filler, respectively, aiming at clarifying key factors such as filler dispersion, filler migration, phase morphology evolution, and their cooperative influences for the formation of thermal conductive networks. The main conclusions are as follows:

In Chapter 2, a reactor granule technology (RGT) was employed to obtain Al<sub>2</sub>O<sub>3</sub> nanoparticles uniformly dispersed in the polypropylene (PP) matrix. It was then mixed with polyolefin elastomer (POE) to obtain a co-continuous stricture. It was found that Al<sub>2</sub>O<sub>3</sub> nanoparticles thermodynamically prefer the localization at the interface between PP and POE, but their successful migration to the interface largely depended on the initial dispersion of the nanoparticles. The agglomerates hindered the migration of fillers to the interface. The uniform

dispersion of Al<sub>2</sub>O<sub>3</sub> nanoparticles achieved by RGT led to the selective localization of Al<sub>2</sub>O<sub>3</sub> nanoparticles at the interface, creating thermal conductive networks along the co-continuous structure, and eventually giving high thermal conductivity. This chapter firstly demonstrated that the distribution of nanoparticles can be controlled even for the combination of hydrophobic and hydrophilic pairs like PP and Al<sub>2</sub>O<sub>3</sub>. In summary, I successfully proposed a facile way to control the localization of filled nanoparticles at the interface in nanocomposites, and key factors thereof in this chapter.

In Chapter 3, I revealed factors that affect the migration of nanoparticles by systematically investigating the nanoparticle migration and phase morphology evolution in the PP/POE/Al<sub>2</sub>O<sub>3</sub> ternary system during the annealing process. By analyzing TEM and SEM micrographs of the nanocomposites at different annealing durations, important cooperation between nanoparticle localization and phase morphology evolution was evidenced: The decreased phase domain size facilitated the successful migration to the interface; Further, the formation of nanoparticle networks at the interface decelerated the phase coarsening during annealing. By employing the RGT, the polymer blend presented a much smaller phase domain size and retarded phase coarsening speed. Therefore, uniformly distributed thermal conductive networks were afforded.

Chapter 4 describes results on the development of thermal conductive PP

nanocomposites based on a continuous segregated structure. The design of the continuous segregated structure aimed at constructing continuous filler networks similar to the selective localization at the interface but with denser thermal conductive paths. By solvent-aided extraction of the POE phase in PP/POE/Al<sub>2</sub>O<sub>3</sub> nanocomposites with a co-continuous structure and subsequent compression of the porous scaffold, the c PP/Al<sub>2</sub>O<sub>3</sub> nanocomposites with continuous segregated distribution gave far greater improvement of thermal conductivity (1.07 W/m K at 27.5 vol%) compared with selective distribution and uncontrolled distribution. The combination of this structure with RGT indicated that both the number and density of thermal conductive paths were important for improving the thermal conductivity.

I believe that the research work carried out in this thesis has established a novel and promising rout for the design of thermal conductive composites in a simple and efficient manner. Furthermore, I have clarified some key factors affecting the formation of thermal conductive networks, which are expected to be useful for other thermal conductive composites design.

### List of publications and other achievements

### **Publications**

1. <u>Xi Zhang</u>, Xiaochao Xia, Hui You, Toru Wada, Patchanee Chammingkwan, Ashutosh Thakur, Toshiaki Taniike, Design of continuous segregated polypropylene/Al<sub>2</sub>O<sub>3</sub> nanocomposites and impact of controlled Al<sub>2</sub>O<sub>3</sub> distribution on thermal conductivity. Composites Part A: Applied Science and Manufacturing 2020; 131: 105825. (IF = 6.28)

2. <u>Xi Zhang</u>, Toru Wada, Patchanee Chammingkwan, Ashutosh Thakur, Toshiaki Taniike, Cooperative influences of nanoparticle localization and phase coarsening on thermal conductivity of polypropylene/polyolefin elastomer blends, Composites Part A: Applied Science and Manufacturing 2019; 126: 105602. (IF = 6.28)

3. <u>Xi Zhang</u>, Bulbul Maira, Yuya Hashimoto, Toru Wada, Patchanee Chammingkwan, Ashutosh Thakur, Toshiaki Taniike, Selective localization of aluminum oxide at interface and its effect on thermal conductivity in polypropylene/polyolefin elastomer blends, Composites Part B: Engineering 2019; 162: 662–70. (IF = 6.86)

### Conferences

### **Domestic Conference**

1. <u>Xi Zhang</u>, Toru Wada, Patchanee Chammingkwan, Toshiaki Taniike, Cooperative Influences of Nanoparticle Localization and Phase Coarsening on Thermal Conductivity of Ternary Nanocomposites, 64th SPSJ Annual Meeting, Fukuoka, Japan, May 27–29, 2020.

### **International Conferences**

1. <u>Xi Zhang</u>, Bulbul Maira, Toru Wada, Patchanee Chammingkwan, Toshiaki Taniike, Selective localization of nanoparticles in immiscible polyolefin blends achieved by reactor granule technology. 5th Blue Sky Conference, Naples and Sorrento, Italy, Jun. 24–28, 2019.

2. <u>Xi Zhang</u>, Bulbul Maira, Toru Wada, Patchanee Chammingkwan, Toshiaki Taniike, Improved thermal conductivity of PP/POE blends by controlling uniform dispersion and selective distribution of Al<sub>2</sub>O<sub>3</sub>, the 12th SPSJ International Polymer Conference (IPC2018), Hiroshima, Japan, Dec. 4–7, 2018.