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Compressibility of the MgB$_2$ superconductor


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Considerable excitement has been caused recently by the discovery that the binary-boride system with stoichiometry MgB$_2$ is superconducting at the remarkably high temperature of 39 K [J. Nagamatsu, N. Nakagawa, T. Muranaka, Y. Zenitani, and J. Akimitsu, Nature 410, 63 (2001)]. This potentially opens the way to even higher-$T_c$ values in a new family of superconductors with unexpectedly simple composition and structure. The simplicity in the electronic and crystal structures could allow the understanding of the physics of high-$T_c$ superconductivity without the presence of the multitude of complicated features, associated with the cuprates. Synchrotron x-ray diffraction was used to measure the isothermal compressibility of MgB$_2$, revealing a stiff tightly packed incompressible solid with only moderate bonding anisotropy between intralayer and interlayer directions. These results, combined with the pressure evolution of the superconducting transition temperature, $T_c$, establish its relation to the B and Mg bonding distances over a broad range of values.

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MgB$_2$ adopts a hexagonal crystal structure (AlB$_2$-type, space group $P6_3/mmm$) that is analogous to intercalated graphite with all hexagonal prismatic sites of the primitive graphic structure (found in hexagonal BN) completely filled and resulting in two interleaved B and Mg layers. In addition, allowing for full charge transfer from Mg to the boron two-dimensional (2D) sheets, the latter are themselves iso-electronic with graphite.

Detailed information on the properties of MgB$_2$ is being, currently, rapidly accumulated. Band structure calculations clearly reveal that, while strong B-B covalent bonding is retained, Mg is ionized, and its two electrons are fully donated to the B-derived conduction band.\(^3\)\(^-\)\(^6\) Superconductivity in MgB$_2$ is then essentially due to the metallic nature of the boron 2D sheets and the presence of strong electron-phonon interactions together with the high-vibrational frequencies of the light B atoms ensure a high-transition temperature.\(^3\) Support for such a phonon-mediated BCS-type mechanism has been provided by measurements of the boron isotope effect ($\Delta T_c = 1.0$ K, isotope exponent $\alpha_B = -0.26$).\(^7\)

In addition, $T_c$ has been found to decrease with applied pressure at the rate of $-dT_c/dP = -1.6$ K/GPa up to 1.84 GPa,\(^8\) again consistent with mediation of the pairing interaction by phonons. An alternative scenario derives from the fact that MgB$_2$ is hole doped and superconductivity may be understood within a formalism developed for high-$T_c$ cuprate superconductivity.\(^9\) Such a theory predicts a positive-pressure coefficient on $T_c$, as a result of the decreasing in-plane B-B distance with increasing pressure and appears to disagree with the experimental observations. However, the response of the system may be more complex if pressure also affects the charge transfer between the B planes and Mg, and will vary depending on whether the system is in the overdoped or underdoped regime.

Here, we address the problem of the evolution of the structural properties of the MgB$_2$ superconductor with applied pressure using synchrotron x-ray powder diffraction techniques. We find that MgB$_2$ remains strictly hexagonal to the highest pressure used. The isothermal interlayer compressibility, $d \ln c/dP$ at zero pressure is only 1.4 times the value of the in-plane compressibility, $d \ln a/dP$, manifesting the anisotropic nature of the crystal structure. Nonetheless, the bonding anisotropy in this material is not as large as other quasi-2D systems like alkali-metal intercalated graphite, and the elastic properties appear only moderately anisotropic. This information is valuable in testing the predictions of competing models for the mechanism of superconductivity. The diffraction experiments in combination with the data on the pressure dependence of $T_c$ permit us to determine its variation over a wide range of unit-cell volume with an initial pressure coefficient, $d \ln T_c/dV = 0.18$ Å$^{-3}$.

The MgB$_2$ sample used in this work was prepared, as reported in Ref. 1, by heating a pressed pellet of stoichiometric amounts of Mg and amorphous B for 10 h at 700 °C under an argon pressure of 196 MPa and was superconducting with $T_c = 39$ K. Phase purity was confirmed by powder x-ray diffraction. The high-pressure synchrotron x-ray diffraction experiments at ambient temperature were performed on beamline BL10XU at Spring-8, Japan. MgB$_2$ was loaded in a diamond anvil cell (DAC), which was used for the high-pressure generation and was equipped with an inconel gasket. The diameter of the top face of the diamond culet was 1 mm, and the sample was introduced in a hole made in the
gasket 0.2 mm deep and 0.4 mm diameter. Silicone oil loaded in the DAC was used as a pressure medium. Pressure was increased at room temperature and was measured with the ruby-fluorescence method. The diffraction patterns were collected using an image-plate detector ($\lambda = 0.49556\ \text{Å}$) with 5 min exposure times. Integration of the two-dimensional diffraction images was performed with the local PIP software and data analysis with the Fullprof suite of Rietveld analysis programs.\(^{10}\) Synchrotron x-ray powder diffraction profiles of MgB$_2$ were collected at pressures between ambient and 6.15 GPa. Inspection of the diffraction data indicated that the pattern could be indexed as hexagonal at all pressures. Thus the refinement of all datasets.\(^{11}\) The Rietveld refinements ($2\theta$ range $= 7 - 43^\circ$) proceeded smoothly (Fig. 1), leading to values for the hexagonal lattice constants, $a = 3.0906(2)\ \text{Å}$ and $c = 3.5287(3)\ \text{Å}$ at ambient pressure (agreement factors: $R_{wp} = 3.9\%$, $R_{B} = 5.9\%$), and $a = 3.0646(1)\ \text{Å}$ and $c = 3.4860(2)\ \text{Å}$ at 6.15 GPa ($R_{wp} = 1.0\%$, $R_{B} = 7.0\%$). Figure 2(a) shows the pressure evolution of the volume of the unit cell of MgB$_2$ together with a least-squares fit of its ambient-temperature equation-of-state (EOS) to the semiempirical second-order Murnaghan EOS:\(^{12}\)

$$P = (K_0/K'_0)[(V_0/V)^{K'_0} - 1],$$

where $K_0$ is the atmospheric-pressure isothermal bulk modulus, $K'_0$ is its pressure derivative ($= dK_0/dP$), and $V_0$ is the unit-cell volume at zero pressure. The fit results in values of $K_0 = 120(5)\ \text{GPa}$ and $K'_0 = 36(3)$. The extracted value of the volume compressibility, $d\ln V/dP = 8.3(3) \times 10^{-3}\ \text{GPa}^{-1}$ implies a stiff tightly packed incompressible solid.

The anisotropy in bonding of the MgB$_2$ structure (Fig. 2 inset) is clearly evident in Fig. 2(b) that displays the variation of the hexagonal lattice constants $a$ and $c$ with pressure. As the applied pressure increases, the $(c/a)$ ratio smoothly decreases (by $\sim 0.4\%$ to 6.15 GPa). We described the pressure dependence of each lattice constant with a variant of Eq. (1), in which $K_0$ and its pressure derivative, $K'_0$ were substituted by the individual $K_x$ and $K'_x (x = a, c)$ values. The results of these fits are also included in Fig. 2(b) and give $K_a = 410(20)\ \text{GPa}$, $K'_a = 13(1)$; $K_c = 292(12)\ \text{GPa}$, $K'_c = 85(7)$. These values clearly reveal the diversity in bonding interactions present, with the solid being least compressible in the basal plane $ab$, in which the covalent B-B bonds lie [d ln $a/dP = 0.0024(1)\ \text{GPa}^{-1}$]. However, the interlayer linear compressibility, $d\ln c/dP = 0.0034(1)\ \text{GPa}^{-1}$ is only 1.4 times larger, implying very stiff Mg-B bonding; significantly, it is considerably smaller than those of the structurally related strongly anisotropic alkali-metal intercalated
TABLE I. Bond distances (Å) for MgB$_2$ at selected pressures.

<table>
<thead>
<tr>
<th>Pressure (GPa)</th>
<th>B-B</th>
<th>Mg-B</th>
<th>Mg-Mg</th>
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<tr>
<td>ambient</td>
<td>1.7844(1)</td>
<td>2.5094(1)</td>
<td>3.0906(2)</td>
</tr>
<tr>
<td>3.23</td>
<td>1.7738(1)</td>
<td>2.4917(1)</td>
<td>3.0724(1)</td>
</tr>
<tr>
<td>6.15</td>
<td>1.7694(1)</td>
<td>2.4837(1)</td>
<td>3.0646(1)</td>
</tr>
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distance and should form a stringent test of competing models for the interpretation of the superconducting pairing mechanism in this material. Detailed band structure calculations should be able to also shed light on the evolution of the charge transfer between Mg and B and decipher the relative importance between the boron layers and the interstitial metal ions to superconductivity.

**Note added in proof**: Recently, we have extended the high-pressure study of MgB$_2$ to 30 GPa using helium and methanol/ethanol as pressure media. No phase transition was identified to this pressure and the bulk moduli have been extracted from Eq. (1) as $K_0=154(8)$ GPa and $K'_0=4.4(8)$.17

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11. The atoms were placed in the following positions in the unit cell: Mg in (1a) (0,0,0); B in (2d) (1/3,2,1/2).


17 Y. Ohishi et al. (unpublished).