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### **CPG based Self-adapting Multi-DOF Robotic Arm Control**

Woosung Yang, Ji-Hun Bae, Yonghwan Oh, Nak Young Chong, Bum-Jae You and Sang-Rok Oh, Member, IEEE

Abstract- Recently, biologically inspired control approaches for robotic systems that involve the use of central pattern generators (CPGs) have been attracting considerable attention owing to the fact that most humans or animals move and walk easily without explicitly controlling their movements. Furthermore, they exhibit natural adaptive motions against unexpected disturbances or environmental changes without considering their kinematic configurations. Inspired by such novel phenomena, this paper endeavors to achieve self-adapting robotic arm motion. For this, biologically inspired CPG based control is proposed. In particular, this approach deals with crucial problems such as motion generation and repeatability of the joints emerged remarkably in most of redundant DOF systems. These problems can be overcome by employing a control based on artificial neural oscillators, virtual force and virtual muscle damping instead of trajectories planning and inverse kinematics. Biologically inspired motions can be attained if the joints of a robotic arm are coupled to neural oscillators and virtual muscles. We experimentally demonstrate self-adaptation motions that that enables a 7-DOF robotic arm to make adaptive changes from the given motion to a compliant motion. In addition, it is verified with real a real robotic arm that human-like movements and motion repeatability are satisfied under kinematic redundancy of joints.

#### I. INTRODUCTION

t is well known that the walking mechanism in the nervous Lsystem by central pattern generators (CPGs) composed of neural oscillators and their network is the fundamental principle for attaining their natural and robust locomotion [1]. Humans or animals exhibit novel natural rhythmic movements such as running, swimming, flying, breathing, etc and continuous arm motions such as turning a steering wheel, rotating a crank, etc. which are dependent upon the interaction between the musculo-skeletal system and the nervous system. Since the musculo-skeletal system is activated like a mechanical spring by means of CPGs and their entrainment property [2]-[5], human behavior is adaptive or robust against unexpected disturbances or environmental changes. Alternate motor commands for the muscles are provided in the CPGs, which enables the musculoskeletal system to deal with environmental perturbations properly afferent feedback of sensory signal

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Relating these previous works, Matsuoka presented a mathematical description of a neural oscillator [2]. He proved that neurons generate a rhythmic patterned output and analyzed the conditions necessary for steady-state oscillations. He also investigated the mutual inhibition networks to control the frequency and pattern [3] but did not include the effect of the feedback on the neural oscillator performance. Employing Matsuoka's neural oscillator model, Taga et al. investigated the sensory signal from the joint angles of a biped robot as feedback signals [4], [5]; they showed that neural oscillators made the robot robust to perturbations through entrainment. Cao et al. [6] proposed the genetic algorithm (GA)-based method to build up desired neural oscillator networks. The CPG based approach was applied later to various locomotion systems [7]-[11] to show that neural oscillators made the robot adaptive to uneven terrains through the entrainment property.

Besides the examples of locomotion, various efforts have been made to strengthen the capability of robots from biological inspiration. Williamson proposed the neuromechanical system that was coupled with the neural oscillator for controlling rhythmic arm motions [12]. Arsenio [13] suggested the multiple-input describing function technique to control multivariable systems connected to multiple neural oscillators. However, they only interest in attaining natural adaptive motions by the coupling between the arm joints and neural oscillators. Thus, the correctness of the desired motion was not guaranteed. Specifically, robot arms are required to exhibit complex behaviors or to trace a trajectory for certain type of tasks, where the substantial difficulty of parameter tuning emerges. Yang et al. presented simulation and experiment results in controlling the robotic arm trajectory incorporating neural oscillators for a desired task [14]-[16].

Apart from such the proposed parameter optimization method, we have addressed an intuitive and efficient approach of biologically inspired control [17]. In addition, this work includes the supplemental method to practically enhance the performance with experimental demonstration using the developed 7-DOF robotic arm. Contributions of this approach can be summarized in the following points. 1) In the CPG based control approach, an imposed task on a multi-DOF robotic arm can be attained easily 2) promoting an impressive capability such as self-adapting motions against an unknown disturbance. 3) Also it is needless to solve ill-posedness problems of inverse kinematics and 4) this approach can give an insight into a method of guaranteeing

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motion repeatability of joints in certain regions of the joint space 5) even in redundancy of degrees of freedom (DOFs). For technically accomplishing these objectives, virtual force constraints in terms of Jacobian transpose and virtual muscle damping factors corresponding to the velocity of joints [14] are employed simply to the CPG based controller as desired torques.

In the following section, the proposed biologically inspired control scheme briefly explained. In Section III, stability of the neural oscillator dynamics is described to design the parameters of the neural oscillator. Details of dynamic responses for the verification of the proposed method through experiments are illustrated and discussed in Section IV. Finally, conclusions are drawn in Section V.

#### II. BIOLOGICALLY INSPIRED ROBOTIC SYSTEMS

# Artificial neural oscillato CPG in the spinal cord Virtual force sculo-skeletal system of human arm and hand

A. Conceptual Model

Fig. 1. Conceptual figure of biologically inspired control for a robotic arm/hand

Muscle

Figure 1 illustrates a schematic model for a robotic arm whose joints are coupled to the neural oscillators. A virtual force leads the coupled robotic arm to a given motion. VFI such as springs and dampers, which are supposed to virtually exist at the target, can be transformed into equivalent torques. This causes the end-effector of a robotic arm to draw according to the target calculating position error. This shows that the ill-posedness of inverse kinematics can be resolved in a natural way without introducing any artificial optimization criterion [18], [19]. However, even with this method, kinematic configurations including redundant joints may not be guaranteed, even though the posture of a robotic arm can only be set within certain boundaries.

From this point of view, it is advantageous if neural oscillators are barely coupled to each joint of a robotic arm. When the oscillators are attached to a robotic arm, they provide proper motor commands that consider the movements of the joints using sensory signals. Since the biologically inspired motions of each joint [17], are attained by the intrinsic entrainment property of the neural oscillator

with its network [20] considering each joint direction, the coupled joint can respond intuitively to environmental changes or unknown disturbances by performing an objective motion. In addition, each neural oscillator can be tuned to produce a criterion in terms of the motion limitation of the joints by considering the amplitude of the sensory feedback signal.

#### B. Artificial Neural Oscillator

Neural motor patterns of vertebrates are obtained from the CPG and modified by sensory signals that detect environmental disturbances. In order to technically accomplish such effects, we used Matsuoka's neural oscillator consisting of two simulated neurons arranged in mutual inhibition, as shown in Fig. 2. If gains are properly tuned, the system exhibits limited cycle behavior. We now propose the control method for dynamic systems that closely interacts with the environment by exploiting the natural dynamics of Matsuoka's oscillator.

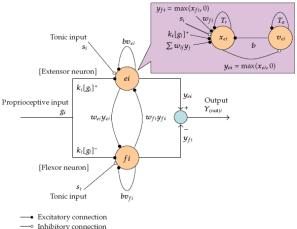


Fig. 2. Schematic diagram of Matsuoka Neural Oscillator

$$T_{r}\dot{x}_{ei} + x_{ei} = -w_{fi}y_{fi} - \sum_{i,j=1}^{n} w_{ij}y_{j} - bv_{ei} - \sum_{i=1}^{n} k_{i}[g_{i}]^{+} + s_{i}$$

$$T_{a}\dot{v}_{ei} + v_{ei} = y_{ei}$$

$$y_{ei} = [x_{ei}]^{+} = \max(x_{ei}, 0), (i = 1, 2, \dots, n)$$

$$T_{r}\dot{x}_{fi} + x_{fi} = -w_{ei}y_{ei} - \sum_{i,j=1}^{n} w_{ij}y_{j} - bv_{fi} - \sum_{i=1}^{n} k_{i}[g_{i}]^{-} + s_{i}$$

$$T_{a}\dot{v}_{fi} + v_{fi} = y_{fi}$$

$$y_{fi} = [x_{fi}]^{+} = \max(x_{fi}, 0), (i = 1, 2, \dots, n)$$
(1)

where  $x_{ei}$  and  $x_{fi}$  indicate the inner state of the *i*th neuron for *i* = 1, 2,..., n, which represents the firing rate. Here, the subscripts e and f denote the extensor and flexor neurons, respectively.  $v_{e(t)i}$  represents the degree of adaptation, and b is the adaptation constant or self-inhibition effect of the *i*th neuron. The output of each neuron  $y_{e(f)i}$  is taken as the positive part of  $x_i$ , and the output of the oscillator is the difference in the output between the extensor and flexor neurons.  $w_{ii}$  is the connecting weight from the *i*th neuron to the *i*th neuron:  $w_{ii}$  is 0 for  $i \neq j$  and 1 for i = j.  $w_{ii}y_i$  represents the total input from neurons arranged to excite one neuron and inhibit the other. These inputs are scaled by the gain  $k_i$ .  $T_r$  and  $T_a$  are the time constants of the inner state and adaptation effect, respectively, and  $s_i$  is an external input with a constant rate.  $w_{e(f)i}$  is the

weight of the extensor neuron or flexor neuron, and  $g_i$  indicates the sensory input from the coupled system that is scaled by the gain  $k_i$ .

#### C. Coupling mechanical systems to neural oscillators

Figure 3 shows a simple mechanical system coupled to the neural oscillator that is simplified by mimicking neuromusculo-skeletal model. The desired torque signal to the *i*-th joint can be given by

$$\tau_i = -k_{oi}(q_i - q_{odi}) - b_i \dot{q}_i, \tag{2}$$

where  $k_{oi}$  is the stiffness of the joint,  $b_i$  the damping coefficient,  $q_i$  the joint angle, and  $q_{odi}$  is the output of the neural oscillator that produces neural commands of the *i*-th joint. The neural oscillator follows the sensory signal from the joints, thus the output of the neural oscillator may change corresponding to the sensory input. This is what is called "entrainment" that can be considered as the tracking of sensory feedback signals so that the mechanical system can exhibit adaptive behavior interacting with the environment.

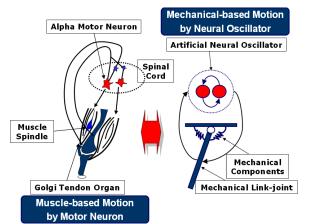


Fig. 3. Neuro-musculo-skeletal model and a simplified biologically inspired mechanical system coupled to the neural oscillator

#### III. STABILITY ANALYSIS OF THE NEURAL OSCILLATOR

All of the neural models suggested have the common feature of neurons being connected so that one neuron's excitation suppresses another's excitation. Matsuoka developed a more general model with the viewpoint that each neuron is capable of receiving different external stimulus and synaptic weights. However, it is difficult to clearly predict the dynamic responses of the artificial neural oscillator incorporated in Matsuoka's work according to specific conditions, where the oscillation, saturation (or convergence), or divergence occur within a certain range. Also, since Matsuoka neural oscillators have nonlinearities such as max(x, x)0) and  $\min(x, 0)$ , as described in (1), it is difficult to analyze their nonlinear behavior. In this section, we discuss the existence of singular points and their stability and use the neural oscillator in time domain analysis to investigate equilibrium states.

Equation (1) of the neural oscillator gives

$$\frac{dx_{e}}{dt} = g_{e}(x_{e}, v_{e}, x_{f}, v_{f}), \frac{dv_{e}}{dt} = f_{e}(x_{e}, v_{e}, x_{f}, v_{f})$$

$$\frac{dx_{f}}{dt} = g_{f}(x_{e}, v_{e}, x_{f}, v_{f}), \frac{dv_{f}}{dt} = f_{f}(x_{e}, v_{e}, x_{f}, v_{f})$$
(3)

where  $g_e, f_e, g_f$ , and  $f_f$  are nonlinear functions of  $x_e, v_e, x_f$ , and  $v_f$ , respectively. Equation (3) can be rewritten as

$$\begin{cases} \dot{x}_{e} \\ \dot{v}_{e} \\ \dot{x}_{f} \\ \dot{v}_{f} \end{cases} = \begin{bmatrix} a_{11} & a_{12} & a_{13} & a_{14} \\ a_{21} & a_{22} & a_{23} & a_{42} \\ a_{31} & a_{32} & a_{33} & a_{43} \\ a_{41} & a_{42} & a_{43} & a_{44} \end{bmatrix} \begin{cases} x_{e} \\ v_{e} \\ x_{f} \\ v_{f} \end{cases}$$
(4)

where

$$\begin{aligned} a_{11} &= \frac{\partial g_e}{\partial x_e} \bigg|_{(0,0,0,0)} & a_{12} &= \frac{\partial g_e}{\partial v_e} \bigg|_{(0,0,0,0)} & a_{13} &= \frac{\partial g_e}{\partial x_f} \bigg|_{(0,0,0,0)} & a_{14} &= \frac{\partial g_e}{\partial v_f} \bigg|_{(0,0,0,0)} \\ a_{21} &= \frac{\partial f_e}{\partial x_e} \bigg|_{(0,0,0,0)} & a_{22} &= \frac{\partial f_e}{\partial v_e} \bigg|_{(0,0,0,0)} & a_{23} &= \frac{\partial f_e}{\partial x_f} \bigg|_{(0,0,0,0)} & a_{24} &= \frac{\partial f_e}{\partial v_f} \bigg|_{(0,0,0,0)} \\ a_{31} &= \frac{\partial g_f}{\partial x_e} \bigg|_{(0,0,0,0)} & a_{32} &= \frac{\partial g_f}{\partial v_e} \bigg|_{(0,0,0,0)} & a_{33} &= \frac{\partial g_f}{\partial x_f} \bigg|_{(0,0,0,0)} & a_{34} &= \frac{\partial g_f}{\partial v_f} \bigg|_{(0,0,0,0)} \\ a_{41} &= \frac{\partial f_f}{\partial x_e} \bigg|_{(0,0,0,0)} & a_{42} &= \frac{\partial f_f}{\partial v_e} \bigg|_{(0,0,0,0)} & a_{43} &= \frac{\partial f_f}{\partial x_f} \bigg|_{(0,0,0,0)} & a_{44} &= \frac{\partial f_f}{\partial v_f} \bigg|_{(0,0,0,0)} \end{aligned}$$

The solutions of (4) should be geometrically similar to those of (3). We assumed the solution of (4) to be in the form

$$\begin{cases} x_e \\ v_e \\ x_f \\ v_f \end{cases} = \begin{cases} X_e \\ V_e \\ X_f \\ V_f \end{cases} e^{\lambda t}$$

$$(5)$$

where  $X_e$ ,  $V_e$ ,  $X_f$ ,  $V_f$ , and  $\lambda$  are constants. Substituting (5) into (4) leads to the eigenvalue problem

$$\begin{bmatrix} -\frac{1}{T_r} -\lambda & -\frac{b}{T_r} & -\frac{w}{T_r} \delta_f & 0\\ \frac{1}{T_a} \delta_e & -\frac{1}{T_a} -\lambda & 0 & 0\\ -\frac{w}{T_r} \delta_e & 0 & -\frac{1}{T_r} -\lambda & -\frac{b}{T_r}\\ 0 & 0 & \frac{1}{T_a} \delta_f & -\frac{1}{T_a} -\lambda \end{bmatrix} \begin{bmatrix} X_e \\ V_e \\ \\ V_e \\ \\ V_f \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$
(6)

The eigenvalues  $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$ , and  $\lambda_4$  can be found by solving the characteristic equation of (6) as

$$\lambda_{1}, \lambda_{2}, \lambda_{3}, \lambda_{4} = \frac{1}{2} (p \pm \sqrt{p^{2} - 4q})$$
 (7)

The stability of the neural oscillator is determined by the nature of the eigenvalues of the state matrix. The various combinations of eigenvalues for the matrix of (6) give various characterizations of the equilibrium of the related nonlinear trajectories in the phase plane. The following then holds: 1) if  $(p^2 - 4q) < 0$ , the motion is oscillatory; 2) if  $(p^2 - 4q) > 0$ , the motion is aperiodic; 3) if p > 0, the system is unstable; and 4) if p < 0, the system is stable. Depending on the eigenvalues in (7), singular or equilibrium points can be classified. Hence, investigating the stability on the inner dynamics of the neural oscillator requires the consideration of possible four conditions:

#### A. Analysis through various cases

1)  $x_e > 0$  and  $x_f > 0$  ( $\delta_e = \delta_f = 1$ )

We can obtain the eigenvalues ( $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$ , and  $\lambda_4$ ) as described in (7)

$$\begin{aligned} \lambda_{1,2} &= -\frac{1}{2T_r T_a} (T_r + (1-w)T_a) \pm \frac{1}{2T_r T_a} \sqrt{(T_r + (1-w)T_a)^2 - 4(b+1-w)T_r T_a} \\ &= \frac{1}{2} \alpha (p_1 \pm \sqrt{(-p_1)^2 - 4q_1}) \end{aligned}$$

$$\begin{aligned} \lambda_{3,4} &= -\frac{1}{2T_r T_a} (T_r + (1+w)T_a) \pm \frac{1}{2T_r T_a} \sqrt{(T_r + (1+w)T_a)^2 - 4(b+1+w)T_r T_a} \\ &= \frac{1}{2} \alpha (p_2 \pm \sqrt{(-p_2)^2 - 4q_2}) \end{aligned}$$

where 
$$p_1 = -(T_r + (1-w)T_a)$$
,  $q_1 = (b+1-w)T_rT_a$ ,  $p_2 = -(T_r + (1+w)T_a)$ ,  $q_2 = (b+1+w)T_rT_a$  and  $\alpha = 1/(T_rT_a)$ .

- Case 1. The eigenvalues  $(\lambda_1, \lambda_2, \lambda_3, \text{ and } \lambda_4)$  are real and distinct  $(p^2 > 4q)$ 

In general, the condition for a root of the equation to have positive values in (8) and (9) is

$$T_a^2 w^2 \pm 2(T_a^2 - T_r T_a) w + T_a^2 + T_r^2 - 2T_a T_r - 4bT_r T_a > 0$$
(10)

To satisfy (10), the following condition needs to hold:

$$\left(T_a - T_r\right)^2 > 4bT_r T_a \tag{11}$$

Under condition (11), the type of motion depends on whether  $\lambda_1, \lambda_2$  and  $\lambda_3, \lambda_4$  are of the same or opposite sign.  $q_1, q_2 > 0$  and  $p_1, p_2 < 0$  such that  $\lambda_1, \lambda_2$  and  $\lambda_3, \lambda_4$  have the same sign. In  $q_1$  and  $q_2 > 0$ ,  $(b+1-w)T_rT_a > 0$  and  $(b+1+w)T_rT_a > 0$ . Then, b > w-1 and b > -w-1 ( $\therefore T_rT_a > 0$ ). The equilibrium when  $\lambda_1 < \lambda_2 < 0$  and  $\lambda_3 < \lambda_4 < 0$  (when  $\lambda_1, \lambda_2$  and  $\lambda_3, \lambda_4$  are real and negative or  $p_1$  and  $p_2 < 0$ ) then becomes stable since  $p_1$  and  $p_2 < 0$ ,  $-(T_r+(1-w)T_a) < 0$ , and  $-(T_r+(1+w)T_a) < 0$ . After rearranging these,  $w < T_r/T_a + 1$ ,  $w > -T_r/T_a - 1$ . Thus, if  $b > \max(-w-1, w-1)$  and  $-T_r/T_a - 1 > w < T_r/T_a + 1$ , the neural oscillator is asymptotically stable and converges to a equilibrium point.

On the other hand, if  $q_1$ ,  $q_2 > 0$  and  $p_1$ ,  $p_2 > 0$ , the eigenvalues have the same positive sign ( $\lambda_1 > \lambda_2 > 0$  and  $\lambda_3 > \lambda_4 > 0$ ), i.e.,  $b > \max(-w-1, w-1)$ ,  $-(T_r+(1-w)T_a) > 0$ , and

 $-(T_r+(1+w)T_a) > 0$ . This condition gives  $w < -T_r/T_a - 1$  and  $T_r/T_a + 1 > w$ . The origin is then at an unstable equilibrium. If  $\lambda_1$ ,  $\lambda_2$  and  $\lambda_3$ ,  $\lambda_4$  are real but of opposite signs ( $q_1$  and  $q_2 < 0$  irrespective of the sign of  $p_1$  and  $p_2$ ), one solution tends to the origin while the other tends to infinity. This is a saddle point with two unstable and two stable manifolds.

Therefore, the stable equilibrium condition can be written as follows:

$$b > \max(-w - 1, w - 1) \text{ and } -T_r/T_a - 1 > w < T_r/T_a + 1$$
 (12)

The unstable equilibrium condition can be written as:

$$b > \max(-w-1, w-1)$$
 and  $w < -T_r/T_a - 1$  and  $w > T_r/T_a + 1(13)$ 

Finally, the condition with a saddle point corresponding to unstable equilibrium is written as:

$$b < w-1 \text{ and } b < -w-1 \tag{14}$$

- Case 2. The eigenvaues  $(\lambda_1, \lambda_2, \lambda_3 \text{ and } \lambda_4)$  are real and equal, respectively.  $(p^2 = 4q)$ 

If  $p_1$  and  $p_2 < 0$  ( $\lambda_1 < 0$  and  $\lambda_2 < 0$ ), the trajectories are straight lines passing through the origin, and the equilibrium points are stable. Otherwise ( $p_1$  and  $p_2 > 0$ ), the origin is unstable. Thus, the stable condition is

$$-T_r/T_a - 1 > w < T_r/T_a + 1 \tag{15}$$

In Case 2, the unstable equilibrium condition can be expressed as

$$w < -T_r/T_a - 1 \text{ and } w > T_r/T_a + 1$$
 (16)

- Case 3. The eigenvaues  $(\lambda_1, \lambda_2, \lambda_3 \text{ and } \lambda_4)$  are complex conjugates  $(p^2 < 4q)$ .

In this case, the stability of motion is determined in terms of the criterion illustrated in Case 1. Hence, Case 3 has the same stable and unstable equilibrium conditions.

2)  $x_e > 0$  and  $x_f < 0$  ( $\delta_e = 1$  and  $\delta_f = 0$ ) or  $x_e < 0$  and  $x_f > 0$  ( $\delta_e = 0$  and  $\delta_f = 1$ )

Both conditions incur similar results. Thus, the eigenvalues are the same if  $x_e$  is exchanged for  $x_f$ . The eigenvalues of (6) in this condition can be obtained as

$$\lambda_1 = -\frac{1}{T_a}, \quad \lambda_2 = -\frac{1}{T_r} \tag{17}$$

$$\lambda_{3,4} = -\frac{1}{2T_r T_a} (T_r + T_a) \pm \frac{1}{2T_r T_a} \sqrt{(T_r + T_a)^2 - 4T_r T_a (b+1)}$$

$$= \frac{1}{2} \alpha (p_2 \pm \sqrt{(-p_2)^2 - 4q_2})$$
(18)

where  $p_2 = -(T_r + T_a)$  and  $q_2 = T_r T_a(b+1)$  and  $\alpha = 1/(T_r T_a)$ .

(9)

- Case 1. The eigenvalues  $(\lambda_3 \text{ and } \lambda_4)$  are real and distinct  $(p^2 > 4q)$ 

In condition (17),  $\lambda_1$  and  $\lambda_2$  always have a negative sign since  $T_r$  and  $T_a > 0$ . These eigenvalues are stable. Therefore, by investigating  $\lambda_3$  and  $\lambda_4$ , the status of the stability can be analyzed. If  $q_2 < 0$ , then the equilibrium is unstable as a saddle point irrespective of the sign of p since  $(b+1)T_rT_a < 0$ , b < -1. In  $q_2 > 0$ , the system's stability is stable. By rearranging these conditions, the stable equilibrium condition is as follows:

$$b > -1$$
 (19)

while the condition for a saddle point corresponding to unstable equilibrium is given as:

$$b < -1 \tag{20}$$

- Case 2. The eigenvaues ( $\lambda_3$  and  $\lambda_4$ ) are real and equal ( $p^2 = 4q$ )

This case satisfies the stable condition since  $p_2 < 0$  and  $\lambda_1$ ,  $\lambda_2$  always have negative signs.

- Case 3. The eigenvaues ( $\lambda_3$  and  $\lambda_4$ ) are complex conjugates ( $p^2 < 4q$ ).

These are stable and unstable equilibrium conditions identical to Case 1.

#### 3) $x_e < 0$ and $x_f < 0$ ( $\delta_e = 0$ and $\delta_f = 0$ )

The eigenvalues  $\lambda_1, \lambda_2, \lambda_3$ , and  $\lambda_4$  can be obtained by solving the characteristic equation of (6) as

$$\lambda_{1,2} = -\frac{1}{T_r}, \ \lambda_{3,4} = -\frac{1}{T_a}$$
(21)

The equilibrium is always stable since all eigenvalues are real and negative. Therefore, all of the trajectories converge to this equilibrium point, and the system does not oscillate.

#### B. Discussion through analysis

If the origin of the equation is unstable or has no stable stationary state, then every solution must be oscillatory (not necessarily periodic) due to the boundedness of the solution. The uniqueness and boundedness of the solution for a neural oscillator was proved by Matsuoka [2], [3] assuming that the total input s from the outside of the network is positive and constant with time. If a root of the equation has a positive real part, the stationary solution is unstable. The basic mutual inhibition network consists of a pair of neurons that reciprocally inhibit each other's excitation.

Thus, with the above cases analyzed in subsection A, oscillations in the neural oscillator are generated by the mutual inhibition condition between n neurons with adaptation. Adaptation plays a key role in oscillation generation [3]. If there was no adaptation effect, oscillation would only occur in networks with special structures. In contrast, if a network has strong adaptation, it can easily generate stable oscillation. When the conditions are satisfied,

the networks produce and sustain oscillation (not necessarily periodic) for any initial state and any temporary disturbance.

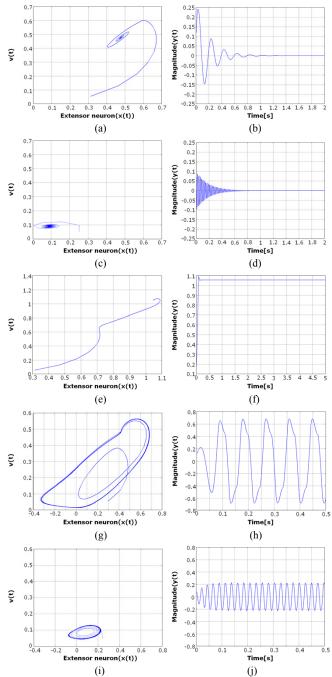


Fig. 4. Plots for analysis of the oscillator dynamics: each row plot is the result when the parameters were sequentially set to w = 10, 10, 10, 15, 100, b = 10, 100, 8.5, 20, 100, and  $T_r/T_a = 10$ , 10, 10, 10, 10, respectively: (a), (c), (e), (g) and (i) phase planes of the extensor neuron (b), (d), (f), (h) and (j) the total outputs of the neural oscillator

In conclusion, based on the analysis of the conditions illustrated in subsection A, the conditions for stable oscillation are

$$w > T_r/T_a + 1 \text{ and } b > w - 1$$
 (22)

Thus, if the final condition of (22) is satisfied and the synaptic weight *w* is large enough (though it must be small compared

with b, the network continues oscillating for any large disturbances. Strictly speaking, the system of the differential equations has no stable equilibrium state for the above situation, only strong mutual inhibition.

Based on the above condition (22), various simulations were carried out to verify the theoretical analysis illustrated in subsection A. As shown in Fig. 4, the analysis was verified as achievable. Solutions for the neural oscillator exist within the asymptotically stable region, according to the results shown in Figs. 4(a) to (d). These results were obtained by employing parameters to incur a Case 1 stable equilibrium state in 1). In the results for Figs 4(a) to (d), the frequency and amplitude of the output of the neural oscillator became higher and smaller with increasing b compared to the result shown in Fig. 4(d) due to the adaptation effect. The saddle point condition induced a stable stationary state for the neural oscillator, as shown in Figs. 4(e) and (f). As seen in Figs. 4(g) to (j), which were obtained from condition (22) in subsection A, the oscillation generated in the neural oscillator was verified to be explicitly present in its unstable condition owing to mutual inhibition between the oscillators. In addition, by comparing the result shown in Figs. 4(g) and (h) with that shown in Figs. 4 (i) and (j) appearances such as the frequency and amplitude exposed from these results is similar at significant points with respect to b, which was revealed when Figure 4(a) and (b) was compared with Fig. 4(c) and (d).

#### IV. VERIFICATION WITH A REAL 7-DOF ROBOTIC ARM

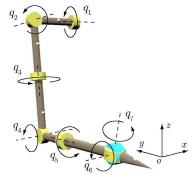


Fig. 5. Geometrical relation of a 7-DOF robotic arm

In this section, we confirm that the neural oscillator and VFI enable the robotic arm to exhibit biologically inspired motion, which enhance its adaptive property sustaining motion stability. Also, in the control scheme, the networks among the oscillators were designed and applied to the proposed approach. The joints q1, q4 and q6 were connected with the excitatory condition of the neural oscillator. There are similar connections to satisfy the condition at the joints q2, q5 and q3, q7, respectively. This is helpful in avoiding the ill-posdness and improving adaptability of the proposed control approach. The proposed control approach (see Appendix) is incorporated to a 7-DOF robotic arm developed by the Korea Institute of Science and Technology (KIST), as seen Figs. 1 and 5. Then, we verify whether or not it is possible to generate a desired movement and adapt to

unknown disturbances while maintaining the repeatability of each joint motion.

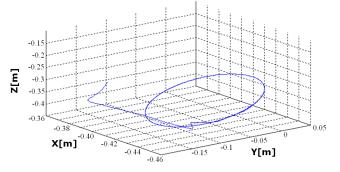


Fig. 6. Trajectory drawn by the end-effector of the 7-DOF robotic arm with respect to the given circular motion in simulation



Fig. 7. Snapshots of the experimental results for the 7-DOF robotic arm with respect to the circular motion when sensory feedback of the neural oscillators was turned off.

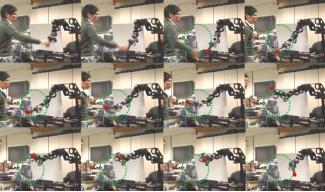


Fig. 8. Snapshots of the experimental results for the 7-DOF robotic arm with respect to the circular motion when sensory feedback of the neural oscillators was turned on

Figure 6 shows the experimental result for the circular motion performed with the real 7-DOF robotic arm. The VFI drives the arm to move according to the given trajectory, and each joint of the robotic arm depends on the outputs generated in the neural oscillator. As expected, the correctness of the arm motion is demonstrated through the result shown in Fig. 6. When controlling a multi-DOF robotic system, if only VFI-based control is considered without using neural oscillators, the repeatability problem for each joint is incurred, as shown in Figs. 7 and 9(a) in contrast with Figs. 8 and 9(b). In Figure 9, the red and blue lines indicate the joint motion

and output of the neural oscillator, respectively. As shown in Fig. 9(b), no alternation of individual joint motions emerged during the circular movement, even in redundant joints, which is in contrast to Fig. 9(a). The joints of the 7-DOF robotic arm were subject to the coupled neural oscillators corresponding to each joint. In Fig. 9(a), the blue line can be ignored because sensory feedback was not fed again. The graph was merely drawn to observe how the neural oscillator was activated. Similar results are shown in Figs. 7 and 8. These figures are the snap shots corresponding to Figs. 9(a) and (b), respectively. Despite the same circular motion being performed in both experiments, the motion in Fig. 7 failed as time went by. In addition, there was an impressive capability for self-adapting motions against unknown disturbance and ill-posedness of inverse kinematics due to the redundant degrees of freedom (DOFs) such as a 7-DOF robotic arm was solved, while the motion repeatability of the joints was sustained.

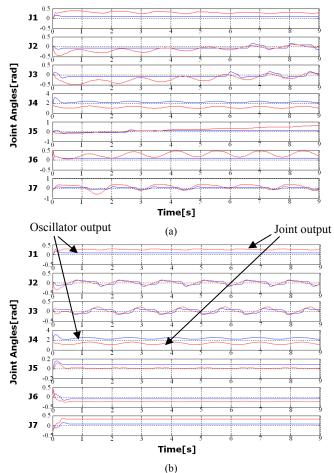


Fig. 9. Individual joint responses acquired during the circular movement in simulation of the 7-DOF robotic arm (a) without coupling to the neural oscillators and (b) with coupling; J1–J7 correspond to each joint in the robotic arm.

Figure 10 illustrates compliant responses that sustained the behaviors carrying out the objectives given to the 7-DOF robotic arm even under unknown disturbances. The red arrow denotes the direction of the applied force. From these results, it is confirmed that the neural oscillator enables the coupled

joint to exhibit a biologically inspired motion to enhance adaptive property sustaining motion stability.

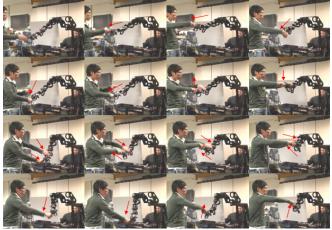


Fig. 10. Snapshots of experimental results for the 7-DOF robotic arm when external disturbances were applied arbitrarily during the circular movement

#### V. CONCLUSION

This work mainly addresses how to achieve human-like behavior of multi-DOF robotic arms employing a biologically inspired control scheme. In order to embody the objective of this work, a neural oscillator and VFI (spring and damper) with a muscle damper are incorporated to the proposed control approach. We first focused on robotic behavior to trace a trajectory correctly with virtual components that make it possible to imposed tasks without considering the ill-posedness even in redundant systems. With this, if the joints of multi-DOF robotic systems are coupled to the neural oscillators as CPGs, the biologically inspired system enables robotic behavior to naturally set kinematic configurations with redundancy and guarantee motion repeatability. The effectiveness of such results was exposed in the experiments on the 7-DOF robotic arm. In addition the robotic arm was demonstrated to adapt in response to environmental changes during the experiments keeping the imposed task.

#### APPENDIX

In general, dynamics of a robot system with *n*-th DOFs could be expressed as

$$H(q)\ddot{q} + \left\{\frac{1}{2}\dot{H}(q) + S(q,\dot{q})\right\}\dot{q} + g(q) = u,$$
(23)

where, *H* denotes the  $n \times n$  inertia matrix of a robot, the second term in the right hand side of (23) stands for coriolis and centrifugal force, and the third term is the gravity effect. Then a control input for a rhythmic motion of the dynamic system shown in (23) is introduced as follows;

$$u = -C_0 \dot{q} - J^T (k\Delta x + \varsigma \sqrt{k} \dot{x}) - k_o \Delta q + g(q), \qquad (24)$$

where

$$C_0 = \operatorname{diag}(c_1, c_2, \cdots, c_n), \ \Delta x = x - x_d, \ \Delta q = q - q_{oa}$$

$$c_i = \zeta_0 \sqrt{k} \sqrt{\sum_{j=1}^n \left[ H_{ij} \right]}, \ (i = 1, \ 2, \ \cdots, \ n)$$

where k and  $\varsigma_0$  is the spring stiffness and damping coefficient, respectively for the virtual components.  $C_0$  is the joint damping.  $k_o$  and  $q_{odi}$  are the stiffness gain and the output of the neural oscillator that produces rhythmic commands, respectively.

The control inputs as seen in (24) consist of two control schemes. One is based on Virtual spring-damper Hypothesis [18], [19] and the other is determined in terms of the output of the neural oscillator as illustrated in (2). In the control input of (24), the first term describes a joint damping for restraining a certain self-motion which could be occurred in a robot system with redundancy, and the second term means PD control in task space by using of Jacobian transpose, and also a spring and a damper in the sense of physics. Appropriate selection of joint damping factors  $C_0$ , stiffness k and damping coefficient  $\varsigma$  render the closed-loop system dynamics convergent, that is, x is converged into  $x_d$  and both of  $\dot{x}$  and  $\dot{q}$  are become 0 as time elapses. In general, the neural oscillators coupled to the joints perform the given motion successively interacting with a virtual constraint owing to the entrainment property, if gains of the neural oscillator are properly tuned [15], [16], [19].

Then, closed-loop dynamics with (23) and (24) is expressed as

$$H(q)\ddot{q} + \left\{\frac{1}{2}\dot{H}(q) + S(q,\dot{q}) + C_0\right\}\dot{q} + J^T(k\Delta x + \varsigma\sqrt{k}\dot{x}) + k_o\Delta q = 0$$
<sup>(25)</sup>

The inner product between  $\dot{q}$  and the closed-loop dynamics of Eq. (25) yields

$$\dot{q}^{T} \left[ H(q)\ddot{q} + \left\{ \frac{1}{2}\dot{H}(q) + S(q,\dot{q}) + C_{0} \right\} \dot{q} + J^{T}k\Delta x + J^{T}\varsigma\sqrt{k}\dot{x} + k_{o}\Delta q \right] = 0^{(26)}$$
And

$$\frac{d}{dt}E = -\dot{q}^T C_0 \dot{q} - \dot{x}^T \varsigma \sqrt{k} \dot{x} \le 0,$$
(27)

where E stands for the total energy

$$E(\dot{q}, \Delta x, \Delta q) = \frac{1}{2} \dot{q}^{T} H(q) \dot{q} + \frac{k}{2} \|\Delta x\|^{2} + \frac{k_{o}}{2} \|\Delta q\|^{2}$$
(28)

In (28), the first term of the quantity *E* describes the kinetic energy of the robot system, the second term means an artificial potential energy caused by the error  $\Delta x$  in task space and the error  $\Delta q$  gives rise to an artificial potential energy corresponding to the third term in joint space. As it is well known in robot control, the energy balance relation of (27) shows that the input-output pair  $(u, \dot{q})$ related to the motion of (26) satisfies passivity.

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